

# Regular DGP in the core of a hypermonopole

Antonio De Felice  
Louvain University  
Beijing – 24 September 2008  
with Christophe Ringeval

**UCL**

# Introduction

- Do we live in  $D > 4$  dimensions?

# Introduction

- Do we live in  $D > 4$  dimensions?
- Deepest idea of '900 – ST motivated

# Introduction

- Do we live in  $D > 4$  dimensions?
- Deepest idea of '900 – ST motivated
- How can gravity be 4D?

# What is gravity?

- ST: KK way, compact small ED

# What is gravity?

- ST: KK way, compact small ED
- RS: non-compact, finite volume ED

# What is gravity?

- ST: KK way, compact small ED
- RS: non-compact, finite volume ED
- DGP: non-compact, infinite volume

- DGP model:

$$S = \frac{M_P^2}{2} \int d^4V R + \frac{M_*^{2+n}}{2} \int d^{n+4}V R$$



- DGP model:

$$S = \frac{M_P^2}{2} \int d^4V R + \frac{M_*^{2+n}}{2} \int d^{n+4}V R$$

- Basis of success and problems [Dvali, Gabadadze, Porrati PLB (2000)][Gregory, Kaloper, Myers, Padilla JHEP (2007)]

- DGP model:

$$S = \frac{M_P^2}{2} \int d^4V R + \frac{M_*^{2+n}}{2} \int d^{n+4}V R$$

- Basis of success and problems [Dvali, Gabadadze, Porrati PLB (2000)][Gregory, Kaloper, Myers, Padilla JHEP (2007)]
- Chosen profile for  $M_*(X)$  [Kolanovic, Porrati, Rombouts PRD (2003)][Shaposnikov, Tinyakov, Zuleta PRD (2004)]

- DGP model:

$$S = \frac{M_P^2}{2} \int d^4V R + \frac{M_*^{2+n}}{2} \int d^{n+4}V R$$

- Basis of success and problems [Dvali, Gabadadze, Porrati PLB (2000)][Gregory, Kaloper, Myers, Padilla JHEP (2007)]
- Chosen profile for  $M_*(X)$  [Kolanovic, Porrati, Rombouts PRD (2003)][Shaposnikov, Tinyakov, Zuleta PRD (2004)]
- Varying mass  $\rightarrow$  ScT-like theory

- DGP model:

$$S = \frac{M_P^2}{2} \int d^4V R + \frac{M_*^{2+n}}{2} \int d^{n+4}V R$$

- Basis of success and problems [Dvali, Gabadadze, Porrati PLB (2000)][Gregory, Kaloper, Myers, Padilla JHEP (2007)]
- Chosen profile for  $M_*(X)$  [Kolanovic, Porrati, Rombouts PRD (2003)][Shaposnikov, Tinyakov, Zuleta PRD (2004)]
- Varying mass  $\rightarrow$  ScT-like theory
- Is there a classical FT which has DGP mechanism?

# Our model

- Topological defects: Classical Field Theory

# Our model

- Topological defects: Classical Field Theory
- 2 ED studied: **no confinement** [Ringeval, Rombouts PRD (2005)]

# Our model

- Topological defects: Classical Field Theory
- 2 ED studied: **no confinement** [Ringeval, Rombouts PRD (2005)]
- DGP can be realized in 7D 'tP-monopole

# Our model

- Topological defects: Classical Field Theory
- 2 ED studied: **no confinement** [Ringeval, Rombouts PRD (2005)]
- DGP can be realized in 7D 'tP-monopole
- 3 ED, associated  $SO(3)$  'tP field configuration



# Our model

- Topological defects: Classical Field Theory
- 2 ED studied: **no confinement** [Ringeval, Rombouts PRD (2005)]
- DGP can be realized in 7D 'tP-monopole
- 3 ED, associated  $SO(3)$  'tP field configuration
- 3 ED asympt flat: min.dim. positively curved

# Action

$$S = \frac{1}{2\kappa^2} \int d^7V e^\psi [R - \partial\psi^2 - U]$$

# Action

$$S = \frac{1}{2\kappa^2} \int d^7V e^\psi [R - \partial\psi^2 - U] \\ + \int d^7V \left[ -\frac{1}{2} D\vec{\Phi}^2 - \frac{1}{4} \vec{H}^2 - \frac{1}{8} \lambda (\vec{\Phi}^2 - v^2)^2 \right]$$

## Action

$$S = \frac{1}{2\kappa^2} \int d^7V e^\psi [R - \partial\psi^2 - U] \\ + \int d^7V \left[ -\frac{1}{2} D\vec{\Phi}^2 - \frac{1}{4} \vec{H}^2 - \frac{1}{8} \lambda (\vec{\Phi}^2 - v^2)^2 \right]$$

- EF,  $U = m_d^2 \psi^2 \exp(\frac{2}{5}\psi)$ ,  $SO(3)$  Higgs  $\vec{\Phi}$

## Action

$$S = \frac{1}{2\kappa^2} \int d^7V e^\psi [R - \partial\psi^2 - U] \\ + \int d^7V \left[ -\frac{1}{2} D\vec{\Phi}^2 - \frac{1}{4} \vec{H}^2 - \frac{1}{8} \lambda (\vec{\Phi}^2 - v^2)^2 \right]$$

- EF,  $U = m_d^2 \psi^2 \exp(\frac{2}{5}\psi)$ ,  $SO(3)$  Higgs  $\vec{\Phi}$
- $D_A \vec{\Phi} = \partial_A \vec{\Phi} - q \vec{C}_A \times \vec{\Phi}$

## Action

$$S = \frac{1}{2\kappa^2} \int d^7V e^\psi [R - \partial\psi^2 - U] \\ + \int d^7V \left[ -\frac{1}{2} D\vec{\Phi}^2 - \frac{1}{4} \vec{H}^2 - \frac{1}{8} \lambda (\vec{\Phi}^2 - v^2)^2 \right]$$

- EF,  $U = m_d^2 \psi^2 \exp(\frac{2}{5}\psi)$ ,  $SO(3)$  Higgs  $\vec{\Phi}$
- $D_A \vec{\Phi} = \partial_A \vec{\Phi} - q \vec{C}_A \times \vec{\Phi}$
- $\vec{H}_{AB} = \partial_A \vec{C}_B - \partial_B \vec{C}_A - q \vec{C}_A \times \vec{C}_B$

# Background geometry

- Static ansatz

$$ds^2 = e^{\sigma(r)} \eta_{\mu\nu} dx^\mu dx^\nu + dr^2 + \omega(r)^2 d\Omega^2$$

# Background geometry

- Static ansatz

$$ds^2 = e^{\sigma(r)} \eta_{\mu\nu} dx^\mu dx^\nu + dr^2 + \omega(r)^2 d\Omega^2$$

- $r, \theta, \phi$  spherical coordinates



# Background geometry

- Static ansatz

$$ds^2 = e^{\sigma(r)} \eta_{\mu\nu} dx^\mu dx^\nu + dr^2 + \omega(r)^2 d\Omega^2$$

- $r, \theta, \phi$  spherical coordinates
- Mapping from  $SO(3)$  to 3 ED's:  $\vec{\Phi} = v f(r) \vec{u}_r$

# Background geometry

- Static ansatz

$$ds^2 = e^{\sigma(r)} \eta_{\mu\nu} dx^\mu dx^\nu + dr^2 + \omega(r)^2 d\Omega^2$$

- $r, \theta, \phi$  spherical coordinates

- Mapping from  $SO(3)$  to 3 ED's:  $\vec{\Phi} = v f(r) \vec{u}_r$

- Gauge fields

$$\vec{C}_\theta = \frac{1 - Q(r)}{q} \vec{u}_\phi \quad \vec{C}_\phi = -\frac{1 - Q(r)}{q} \sin \theta \vec{u}_\theta$$

# Boundary conditions

- $f \rightarrow 1, Q \rightarrow 0$  as  $r \rightarrow \infty$

## Boundary conditions

- $f \rightarrow 1, Q \rightarrow 0$  as  $r \rightarrow \infty$
- $f(0) = 0, Q(0) = 1$

## Boundary conditions

- $f \rightarrow 1, Q \rightarrow 0$  as  $r \rightarrow \infty$
- $f(0) = 0, Q(0) = 1$
- $\sigma = 0, \omega \rightarrow r, \psi \rightarrow 0$  as  $r \rightarrow \infty$

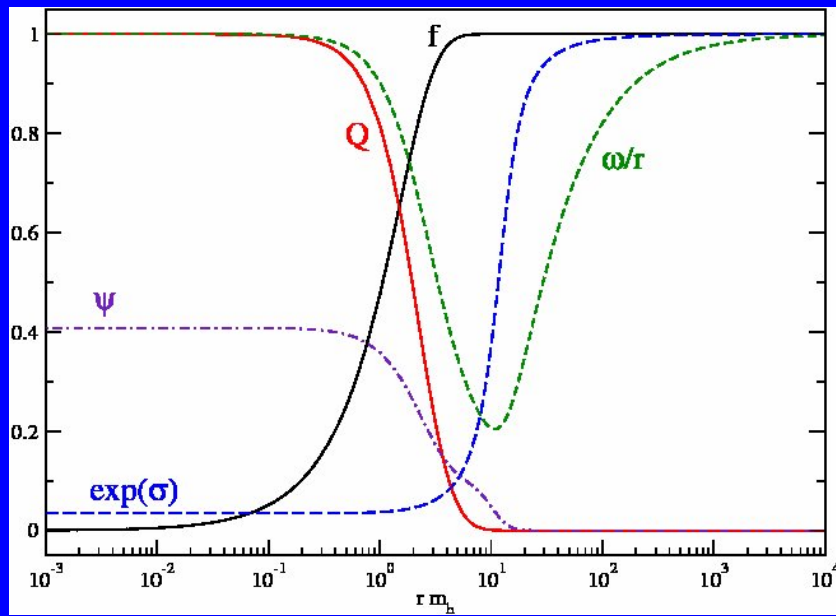
## Boundary conditions

- $f \rightarrow 1, Q \rightarrow 0$  as  $r \rightarrow \infty$
- $f(0) = 0, Q(0) = 1$
- $\sigma = 0, \omega \rightarrow r, \psi \rightarrow 0$  as  $r \rightarrow \infty$
- $\sigma'(0) = \psi'(0) = 0, \omega \sim r$  as  $r \rightarrow 0$

## Boundary conditions

- $f \rightarrow 1, Q \rightarrow 0$  as  $r \rightarrow \infty$
- $f(0) = 0, Q(0) = 1$
- $\sigma = 0, \omega \rightarrow r, \psi \rightarrow 0$  as  $r \rightarrow \infty$
- $\sigma'(0) = \psi'(0) = 0, \omega \sim r$  as  $r \rightarrow 0$
- Define

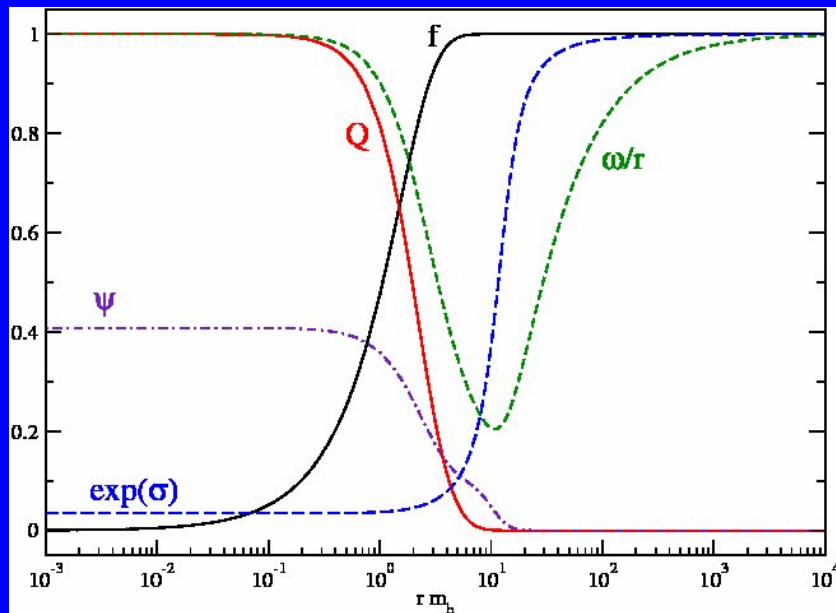
$$\alpha \equiv \kappa^2 v^2 \quad \epsilon \equiv \frac{q^2 v^2}{\lambda v^2} \quad \beta \equiv \frac{m_d^2}{\lambda v^2}$$



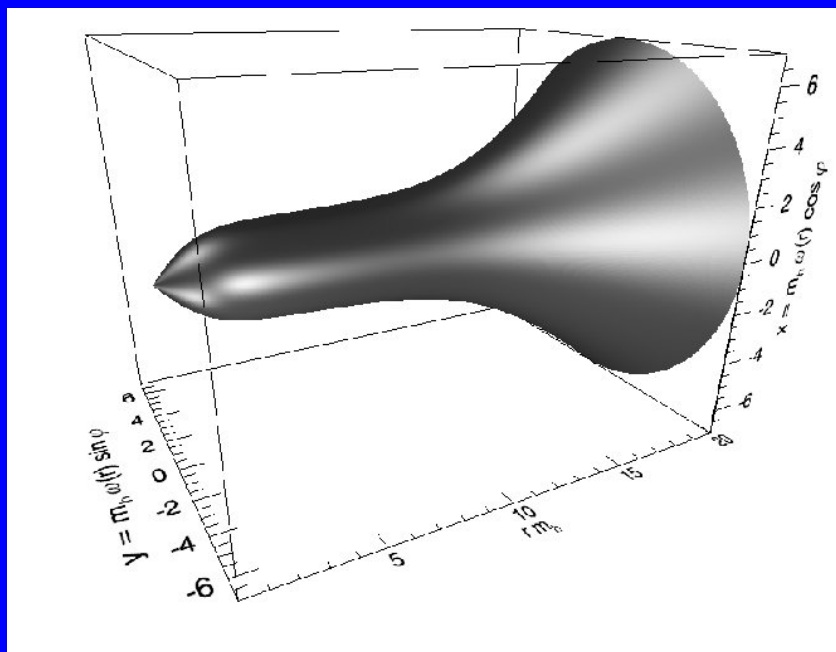
9

Example  $\alpha = 2$ ,  $\epsilon = 0.5$ ,  
 $\beta = 1$ . Typical topological  
 configuration for the fields.  
 Minkowski at boundary.  
 TWPBVP: Cash and  
 Mazza algorithm, JCAM  
 (2005).





Example  $\alpha = 2$ ,  $\epsilon = 0.5$ ,  $\beta = 1$ . Typical topological configuration for the fields. Minkowski at boundary. TWPBVP: Cash and Mazza algorithm, JCAM (2005).



Region where  $\omega$  no longer grows, cylindrically shaped ED, resonant gravitons at this scale

- Two methods of integrations:

- Two methods of integrations:
- – Shooting from boundaries to middle point

- Two methods of integrations:
- – Shooting from boundaries to middle point
- – Adv. Relax. TWPBVP [Cash-Mazzia, 2005]

- Two methods of integrations:
- – Shooting from boundaries to middle point
- – Adv. Relax. TWPBVP [Cash-Mazzia, 2005]
- Same results

- Two methods of integrations:
  - – Shooting from boundaries to middle point
  - – Adv. Relax. TWPBVP [Cash-Mazzia, 2005]
- Same results
- Monopole: no singularities

# Tensor fluctuations

- Introduce TT pert:  $\eta_{\mu\nu} \rightarrow \eta_{\mu\nu} + h_{\mu\nu}$

# Tensor fluctuations

- Introduce TT pert:  $\eta_{\mu\nu} \rightarrow \eta_{\mu\nu} + h_{\mu\nu}$
- $z \equiv m_h \int e^{-\sigma/2} dr$ ,  $\xi_{\mu\nu} \equiv e^{\psi/2+3\sigma/4} \omega h_{\mu\nu}$



# Tensor fluctuations

- Introduce TT pert:  $\eta_{\mu\nu} \rightarrow \eta_{\mu\nu} + h_{\mu\nu}$
- $z \equiv m_h \int e^{-\sigma/2} dr$ ,  $\xi_{\mu\nu} \equiv e^{\psi/2+3\sigma/4} \omega h_{\mu\nu}$
- Linear PDE

$$-\frac{\partial^2 \xi}{\partial z^2} + \left( W^2 + W' - \frac{e^\sigma}{m_h^2 \omega^2} L^2 - \square \right) \xi = 0$$

## Tensor fluctuations

- Introduce TT pert:  $\eta_{\mu\nu} \rightarrow \eta_{\mu\nu} + h_{\mu\nu}$
- $z \equiv m_h \int e^{-\sigma/2} dr$ ,  $\xi_{\mu\nu} \equiv e^{\psi/2+3\sigma/4} \omega h_{\mu\nu}$
- Linear PDE

$$-\frac{\partial^2 \xi}{\partial z^2} + \left( W^2 + W' - \frac{e^\sigma}{m_h^2 \omega^2} L^2 - \square \right) \xi = 0$$

- FT+SH:  $\square \rightarrow M^2$  and  $L^2 \rightarrow -l(l+1)$ ,  $\xi(0) \rightarrow 0$ .

# SUSY-QM Schroedinger equation

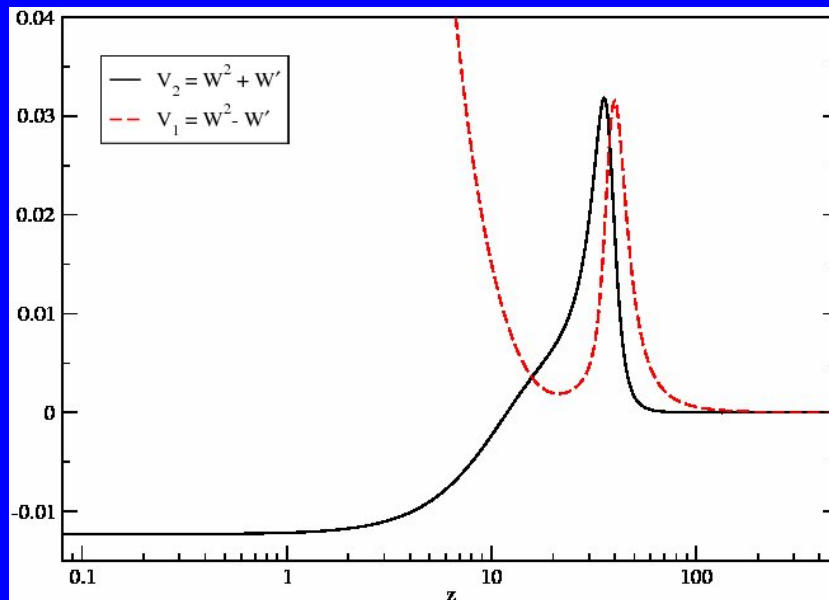
- $W = \frac{3}{4}\sigma' + \omega'/\omega + \frac{1}{2}\psi'$

# SUSY-QM Schroedinger equation

- $W = \frac{3}{4}\sigma' + \omega'/\omega + \frac{1}{2}\psi'$
- QM with central potential  $V_2 \equiv W^2 + W'$

# SUSY-QM Schroedinger equation

- $W = \frac{3}{4}\sigma' + \omega'/\omega + \frac{1}{2}\psi'$
- QM with central potential  $V_2 \equiv W^2 + W'$



Depends on  $\alpha, \beta, \epsilon$ . For  $V_2$ : ground state  $\xi_0$  not normalizable. For  $V_1$ ,  $1/\xi_0$  not regular in 0. SUSY is broken and spectrum  $M^2 > 0$ .

- Orthonormal basis  $u_{M,l}(z)$  solution for Schr. eq.

$$\int_0^{\infty} u_{M,l}^*(z_1) u_{M,l}(z_2) dM = \delta(z_1 - z_2)$$

- Orthonormal basis  $u_{M,l}(z)$  solution for Schr. eq.

$$\int_0^\infty u_{M,l}^*(z_1) u_{M,l}(z_2) dM = \delta(z_1 - z_2)$$

- Using retarded Green function, for TT source  $S_{\mu\nu}$  [Garriga,

Tanaka PRL (2000)]

$$h_{\mu\nu}(X_1) = -\frac{2\kappa^2}{m_h^2 \omega} e^{-\psi(z_1) - 3\sigma/4} \\ \times \int G_\xi(X_1; X_2) e^{-\psi(z_2) + 3\sigma/4} \omega(z_2) S_{\mu\nu} d^7 X_2$$

## Flat case

- Setting  $\psi = \sigma = 0, \omega = r$



## Flat case

- Setting  $\psi = \sigma = 0, \omega = r$
- Source  $S_{\mu\nu} = z^{-2}\delta(z)\delta(\cos\theta)\delta(\phi)s_{\mu\nu}$

## Flat case

- Setting  $\psi = \sigma = 0, \omega = r$
- Source  $S_{\mu\nu} = z^{-2}\delta(z)\delta(\cos\theta)\delta(\phi)s_{\mu\nu}$
- Radial solution:  $u_{M,l}^b = \sqrt{Mz}J_{l+1/2}(Mz)$

## Flat case

- Setting  $\psi = \sigma = 0, \omega = r$
- Source  $S_{\mu\nu} = z^{-2}\delta(z)\delta(\cos\theta)\delta(\phi)s_{\mu\nu}$
- Radial solution:  $u_{M,l}^b = \sqrt{Mz}J_{l+1/2}(Mz)$
- Only  $l = 0$  mode gives contribution

## Flat case

- Setting  $\psi = \sigma = 0, \omega = r$
- Source  $S_{\mu\nu} = z^{-2}\delta(z)\delta(\cos\theta)\delta(\phi)s_{\mu\nu}$
- Radial solution:  $u_{M,l}^b = \sqrt{Mz}J_{l+1/2}(Mz)$
- Only  $l = 0$  mode gives contribution
- Solution

$$h_{\mu\nu}^b = \lim_{z \rightarrow 0} \frac{\kappa^2}{8\pi^2 m_h^2} \int d^3\vec{x}_2 s_{\mu\nu} \int dM \frac{|u_{M,0}^b(z)|^2 e^{-M\Delta\vec{x}}}{z^2 |\Delta\vec{x}|}$$

- Bessel expansion  $u_{M,0}^b \sim \sqrt{2/\pi M} z$

- Bessel expansion  $u_{M,0}^b \sim \sqrt{2/\pi} M z$
- Solution simplifies to

$$h_{\mu\nu}^b = \frac{2\kappa^2}{4\pi^3 m_h^2} \int d^3 \vec{x}_2 \frac{s_{\mu\nu}(\vec{x}_2)}{|\Delta \vec{x}|^4}$$

- Bessel expansion  $u_{M,0}^b \sim \sqrt{2/\pi} M z$
- Solution simplifies to

$$h_{\mu\nu}^b = \frac{2\kappa^2}{4\pi^3 m_h^2} \int d^3 \vec{x}_2 \frac{s_{\mu\nu}(\vec{x}_2)}{|\Delta \vec{x}|^4}$$

- Power law dependence  $1/|\Delta \vec{x}|^{d-2}$  in  $d = 6$  spatial dimensions

- Bessel expansion  $u_{M,0}^b \sim \sqrt{2/\pi} M z$
- Solution simplifies to

$$h_{\mu\nu}^b = \frac{2\kappa^2}{4\pi^3 m_h^2} \int d^3 \vec{x}_2 \frac{s_{\mu\nu}(\vec{x}_2)}{|\Delta \vec{x}|^4}$$

- Power law dependence  $1/|\Delta \vec{x}|^{d-2}$  in  $d = 6$  spatial dimensions
- $4\pi^3$ :  $d - 2$  times the surface of  $d - 1$  unit sphere.



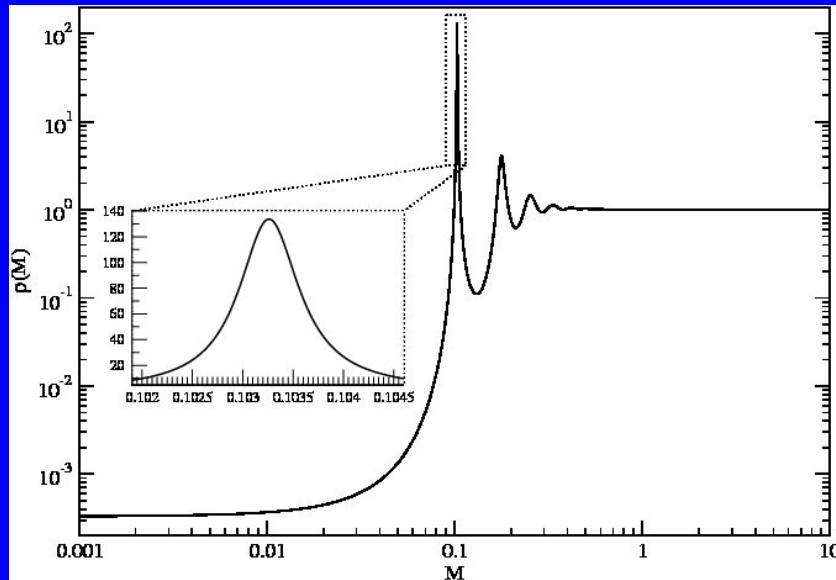
# Monopole case

- Need spectral density  $\rho(M) \equiv |u_{M,0}(0)|^2 / |u_{M,0}^b(0)|^2$

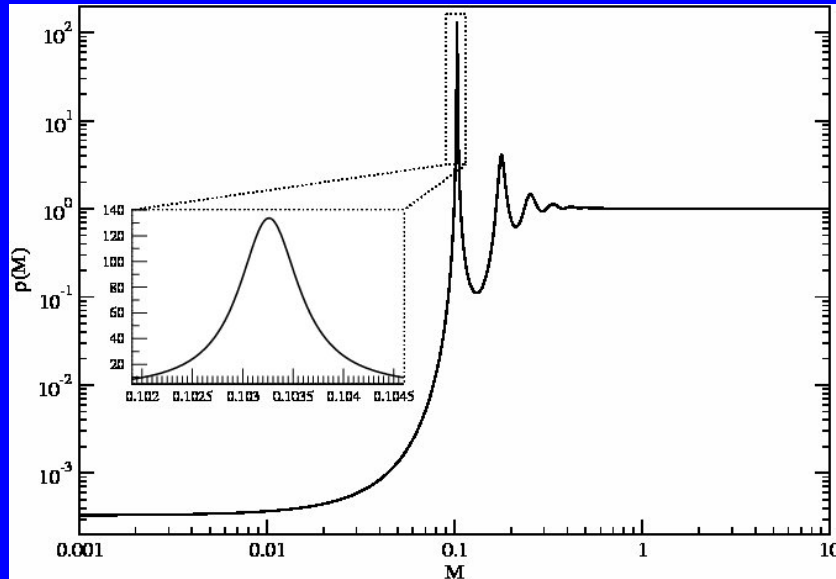
# Monopole case

- Need spectral density  $\rho(M) \equiv |u_{M,0}(0)|^2 / |u_{M,0}^b(0)|^2$
- Therefore one finds

$$h_{\mu\nu} = \frac{2\kappa^2}{8\pi^3 m_h^2} \int d^3 \vec{x}_2 \frac{\mathcal{L}\{\rho(M) M^2\}}{|\Delta \vec{x}|} s_{\mu\nu}(\vec{x}_2)$$

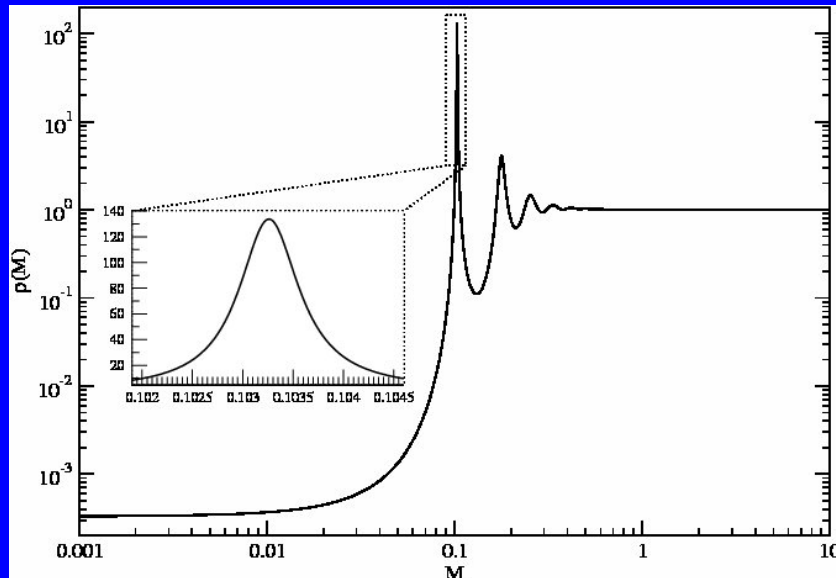


There are at least two trapped **metastable gravitons**. Mass  $m_g \approx 0.103$  Breit-Wigner with  $\Gamma \approx 3.5 \times 10^{-4}$  ( $m_h$  units)



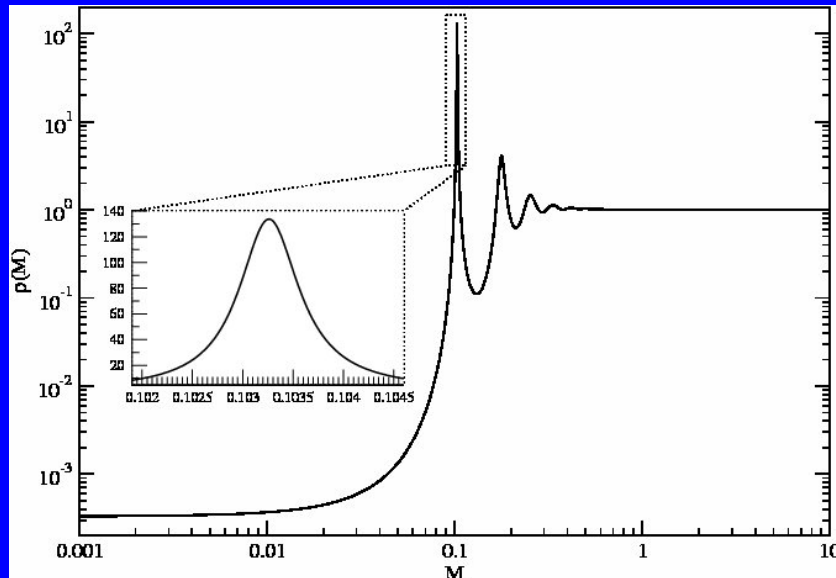
There are at least two trapped **metastable gravitons**. Mass  $m_g \approx 0.103$  Breit-Wigner with  $\Gamma \approx 3.5 \times 10^{-4}$  ( $m_h$  units)

- Constant  $\rho$ : 7D gravity



There are at least two trapped **metastable gravitons**. Mass  $m_g \approx 0.103$  Breit-Wigner with  $\Gamma \approx 3.5 \times 10^{-4}$  ( $m_h$  units)

- Constant  $\rho$ : 7D gravity
- $\rho$  strongly peaked: **resonant metastable modes**



There are at least two trapped **metastable gravitons**. Mass  $m_g \approx 0.103$  Breit-Wigner with  $\Gamma \approx 3.5 \times 10^{-4}$  ( $m_h$  units)

- Constant  $\rho$ : 7D gravity
- $\rho$  strongly peaked: **resonant metastable modes**
- No  $M^2 < 0$  bound state

# DGP mechanism

- Trapping a graviton,  $\rho \sim 1 + C\delta(M - m_g)$

## DGP mechanism

- Trapping a graviton,  $\rho \sim 1 + C\delta(M - m_g)$
- $\mathcal{L}\{\rho M^2\} = 2/|\Delta\vec{x}|^3 + Cm_g^2 e^{-m_g|\Delta\vec{x}|}$



## DGP mechanism

- Trapping a graviton,  $\rho \sim 1 + C\delta(M - m_g)$
- $\mathcal{L}\{\rho M^2\} = 2/|\Delta\vec{x}|^3 + Cm_g^2 e^{-m_g|\Delta\vec{x}|}$
- 4D gravity:  $(m_g/C)^{1/3} < |\Delta\vec{x}|m_g < 1$ , and  $m_g < C$

# DGP mechanism

- Trapping a graviton,  $\rho \sim 1 + C\delta(M - m_g)$
- $\mathcal{L}\{\rho M^2\} = 2/|\Delta\vec{x}|^3 + Cm_g^2 e^{-m_g|\Delta\vec{x}|}$
- 4D gravity:  $(m_g/C)^{1/3} < |\Delta\vec{x}|m_g < 1$ , and  $m_g < C$
- 7D at small/large distance

# DGP mechanism

- Trapping a graviton,  $\rho \sim 1 + C\delta(M - m_g)$
- $\mathcal{L}\{\rho M^2\} = 2/|\Delta\vec{x}|^3 + Cm_g^2 e^{-m_g|\Delta\vec{x}|}$
- 4D gravity:  $(m_g/C)^{1/3} < |\Delta\vec{x}|m_g < 1$ , and  $m_g < C$
- 7D at small/large distance
- We observed fractional power

## Green function

$$G_\xi'' + (-V_2 + \square + e^\sigma \omega^{-2} L) G_\xi = \delta^4(x^\mu - x^{\mu'}) \\ \times \delta(\cos \theta - \cos \theta') \delta(\phi - \phi') \delta(z - z')$$

$$G_\xi = - \int \frac{d^4 p}{(2\pi)^4} e^{ip_\mu(x_1^\mu - x_2^\mu)} \sum_{l,m} Y_l^m(\theta_1, \phi_1) Y_l^{m*}(\theta_2, \phi_2) \\ \times \int \frac{u_{M,l}(z_1) u_{M,l}^*(z_2) dM}{M^2 + \vec{p}^2 - (p^0 + i\epsilon)^2}$$