

Dynamical models for dark matter halos with universal properties

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Dark matter halos

N -body simulations have become the standard way to investigate the structure, dynamics and evolution of dark matter halos.

They revealed several “universal” properties:

- A “universal” density profile
- A power-law pseudo-phase space density
- A linear density slope – velocity dispersion relation

$$\gamma(r) = -\frac{d \ln \rho}{d \ln r}(r) = \frac{\gamma_0 + \gamma_\infty (r/r_s)^\eta}{1 + (r/r_s)^\eta}$$

$$Q(r) = \rho / \sigma^3(r) \propto r^{-\alpha}$$

$$\beta(\gamma) \simeq 1 - 1.15(1 + \gamma/6)$$

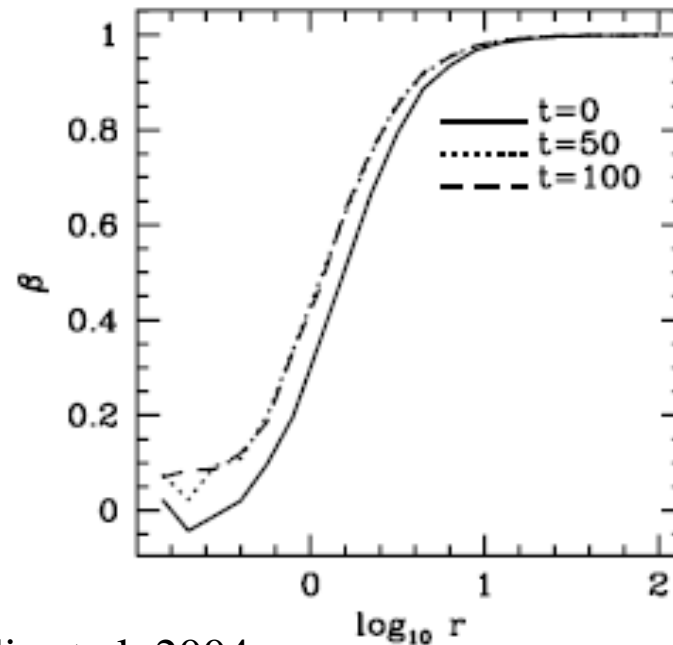
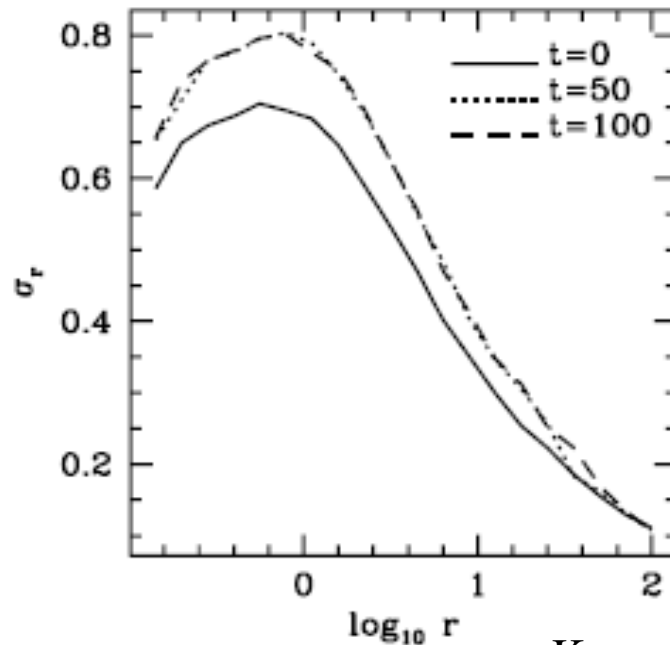
$$\beta(r) = 1 - \frac{\sigma_\theta^2(r)}{\sigma_r^2(r)}$$

Can we construct analytical dynamical models with these properties?

Analytical dynamical models

Equilibrium halos described by a phase-space DF $F(\vec{r}, \vec{v})$

- “noise-free” description of the dynamical structure
 - simple toy models to represent a galaxy or dark matter halo
 - get a better insight into the phase-space structure
- Generate the initial conditions for realistic N-body models

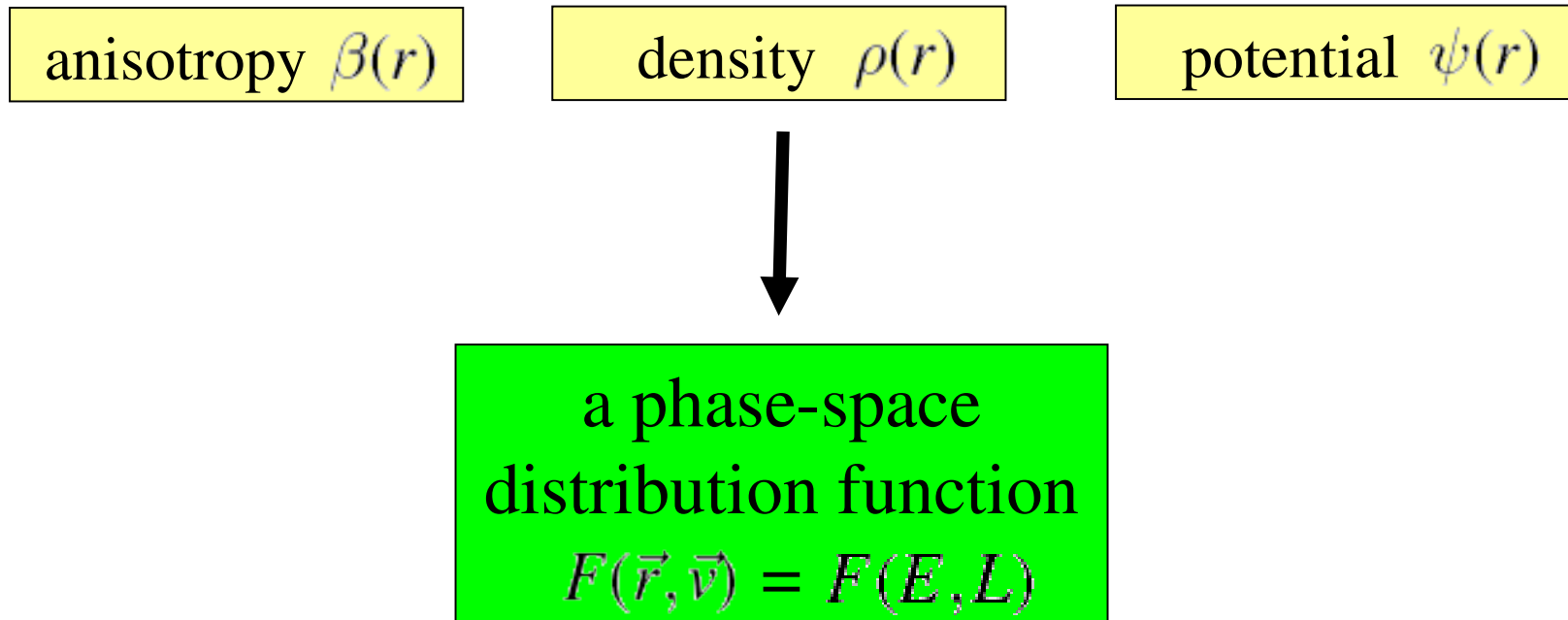


Constructing dynamical models

Dynamical model = phase-space distribution function $F(\vec{r}, \vec{v})$

- contains all possible dynamical information
- can be written in terms of the integrals of motion
- yields direct view on the orbital structure

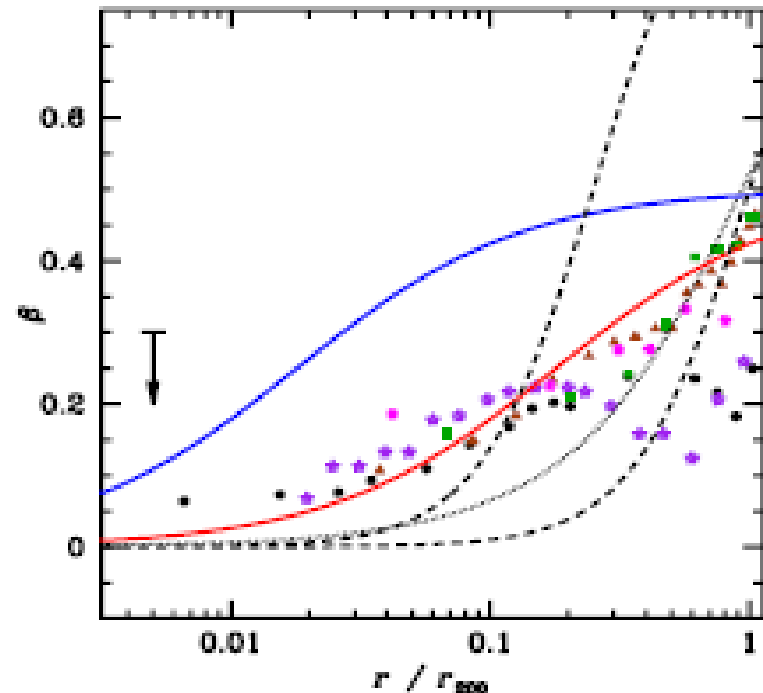
spherical case



Constructing dynamical models

- mathematics are extremely complicated (non-linear integral transforms)
- only analytical solutions for a few sets of anisotropy constraints
 - systems with isotropy or a constant anisotropy
 - Osipkov-Merritt and Cuddeford models (completely radially anisotropic at large radii)

These assumptions are not realistic



More realistic models...

- how to build more realistic models ?
- **augmented density** formalism:
one-to-one relation between $F(E, L)$ and $\tilde{\rho}(\psi, r)$
- problem: formulae complicated and numerically unstable...

$$\tilde{\rho}(\psi, r) = 2\pi M_{\text{tot}} \int_0^\psi dE \int_0^{2(\psi-E)} \frac{F(E, r v_T)}{\sqrt{2(\psi-E) - v_T^2}} dv_T^2$$

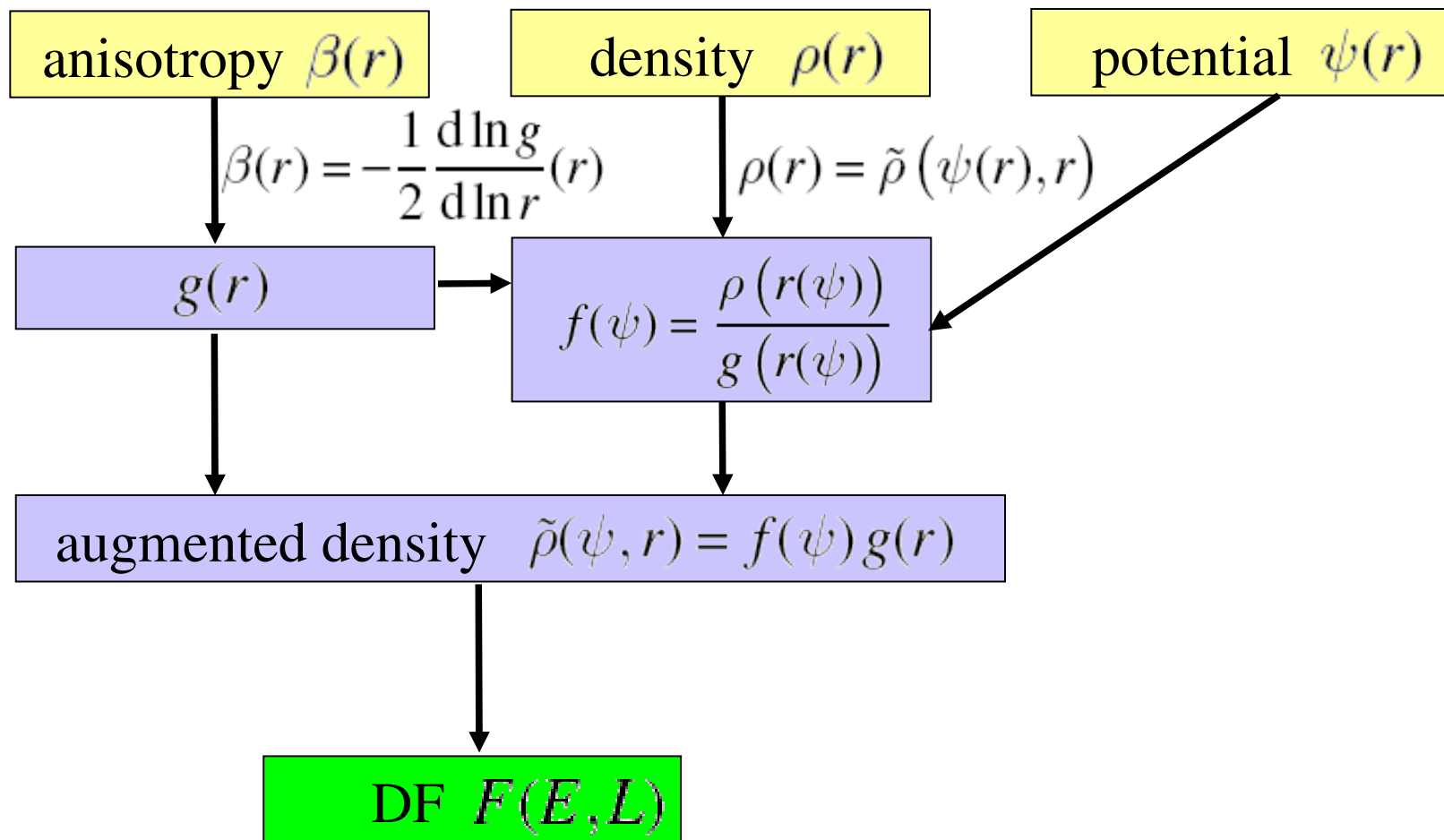
Dejonghe 1986

- things are easier for separable functions

$$\tilde{\rho}(\psi, r) = f(\psi) g(r)$$

the anisotropy is completely determined by $g(r)$

More realistic models...



Practical construction

derivation of the augmented density looks fine in theory....
practically impossible for general anisotropy profiles

⇒ solution: parameterized functions

$$\beta(r) = \frac{\beta_0 + \beta_\infty (r/r_a)^{2\delta}}{1 + (r/r_a)^{2\delta}}$$
$$g(r) = \left(\frac{r}{r_a}\right)^{-2\beta_0} \left(1 + \frac{r^{2\delta}}{r_a^{2\delta}}\right)^{\beta_\infty}$$
$$f_i(\psi) = \rho_{0i} \left(\frac{\psi}{\psi_0}\right)^{p_i} \left(1 - \frac{\psi^{s_i}}{\psi_0^{s_i}}\right)^{q_i}$$

the anisotropy profile is a strong generalization of the special cases mentioned before
(constant anisotropy, Osipkov-Merritt,...)

Analytical models

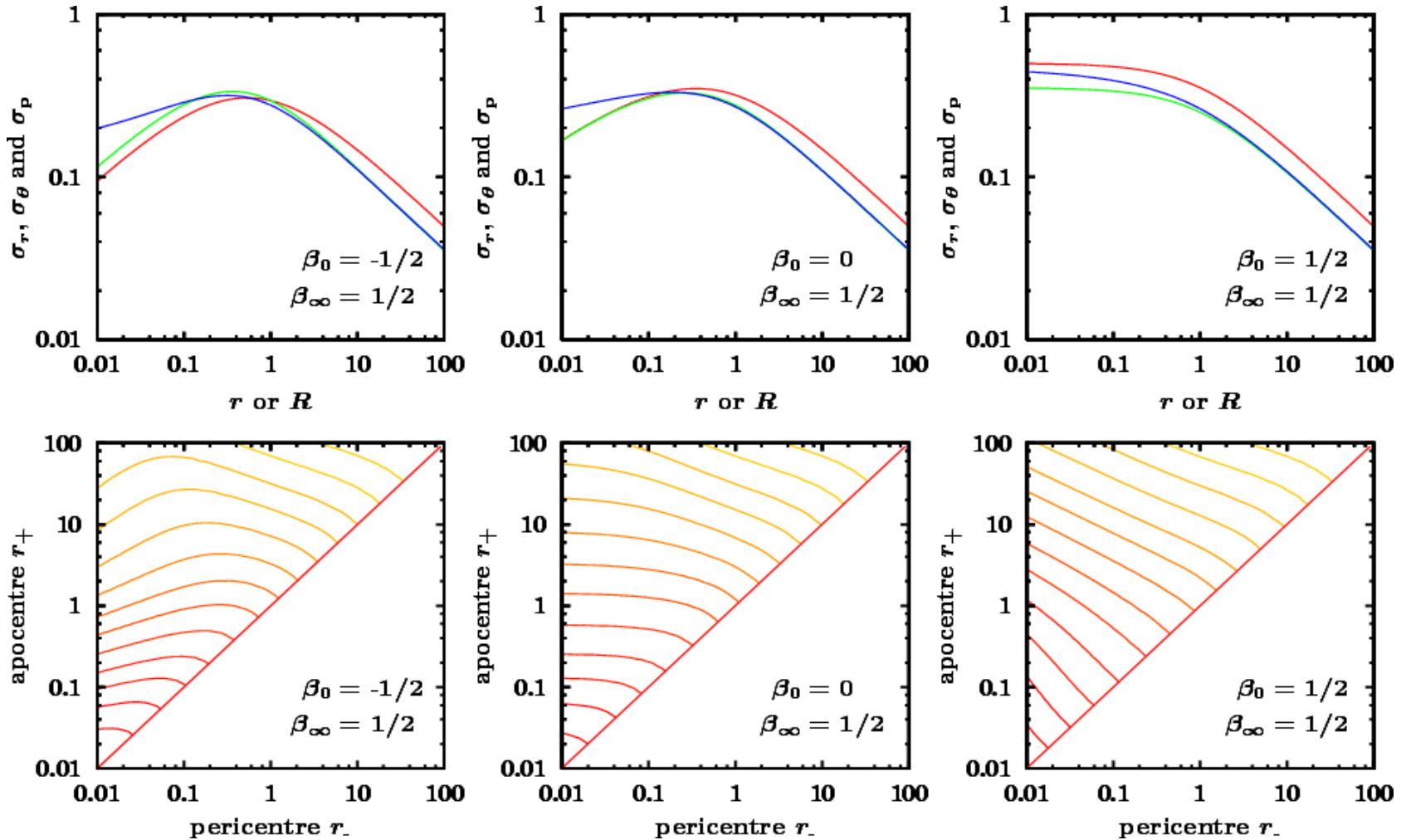
- a general seven-parameter family of augmented densities

$$\tilde{\rho}_i(\psi, r) = \rho_{0i} \left(\frac{\psi}{\psi_0} \right)^{p_i} \left(1 - \frac{\psi^{s_i}}{\psi_0^{s_i}} \right)^{q_i} \left(\frac{r}{r_a} \right)^{-2\beta_0} \left(1 + \frac{r^{2\delta}}{r_a^{2\delta}} \right)^{\beta_\delta}$$

$$F_i(E, L) = \frac{\rho_{0i}}{M_{\text{tot}}(2\pi\psi_0)^{3/2}} \sum_{j=0}^{\infty} (-1)^j \binom{q_i}{j} \left(\frac{E}{\psi_0} \right)^{p_i + js_i - 3/2} \\ \times \sum_{k=0}^{\infty} \binom{\beta_\delta}{k} \frac{\Gamma(1 + p_i + js_i)}{\Gamma(p_i + js_i - \frac{1}{2} + \beta_k) \Gamma(1 - \beta_k)} \left(\frac{L^2}{2r_a^2 E} \right)^{-\beta_k}$$

- contains (a.o.) self-consistent Plummer and Hernquist models
 - with arbitrary anisotropy at small and large radii
 - completely analytical dispersions, DF,...

Analytical Hernquist models



Baes & Van Hese 2007, A&A, 471, 419

how to generate arbitrary densities?

More general analytical models

Quadratic Programming (Dejonghe 1989):

- construct a library of different components with the same anisotropy

$$\tilde{\rho}_i(\psi, r) = f_i(\psi) \left(\frac{r}{r_a} \right)^{-2\beta_0} \left(1 + \frac{r^{2\delta}}{r_a^{2\delta}} \right)^{\beta_\delta}$$

- create a set of M density data points $\rho_{\text{obs}}(r_m)$
- calculate the corresponding $\rho_i(r_m) = \tilde{\rho}_i(\psi(r_m), r_m)$

- select the best-fitting linear combination of N components to the density data

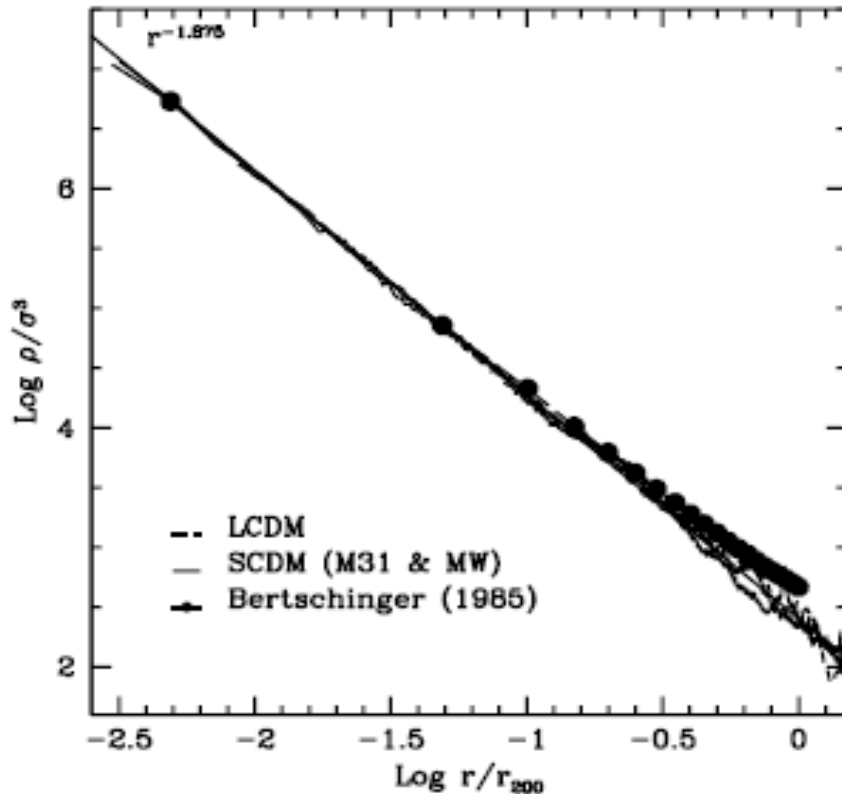
$$\chi_N^2 = \frac{1}{M} \sum_{m=1}^M w_m \left(\rho_{\text{obs}}(r_m) - \sum_{i=1}^N a_{N,i} \rho_i(r_m) \right)^2$$

- the corresponding DF is then simply

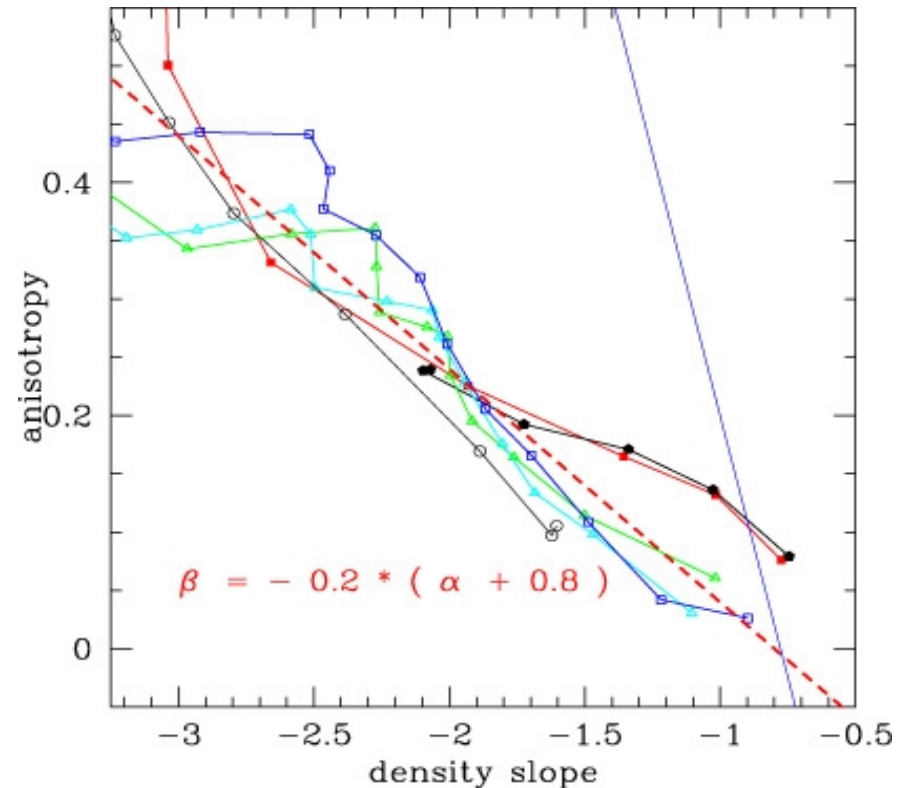
$$F(E, L) = \sum_{i=1}^N a_{N,i} F_i(E, L)$$

Empirical properties of dark halos

- typically cusped in the central regions
- power-law behavior for “pseudo” phase-space density
- linear relationship between density slope and anisotropy



Taylor & Navarro 2001



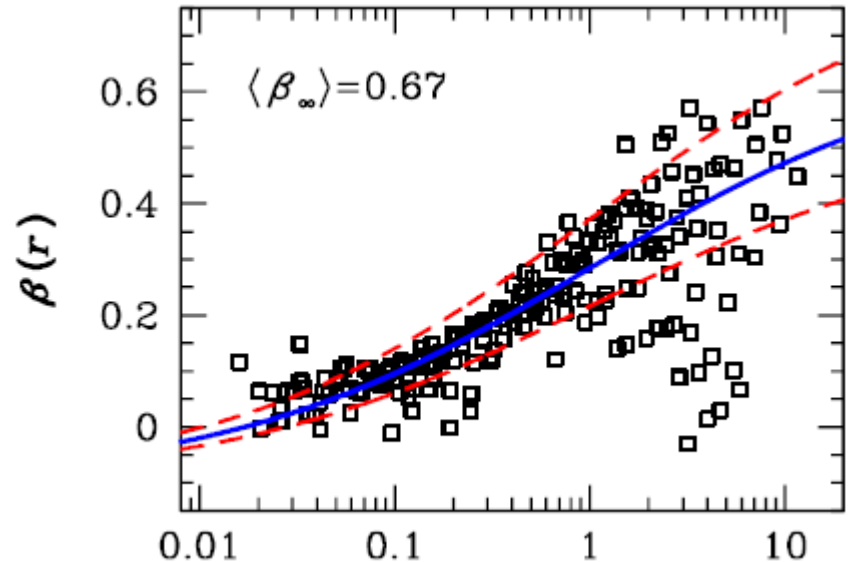
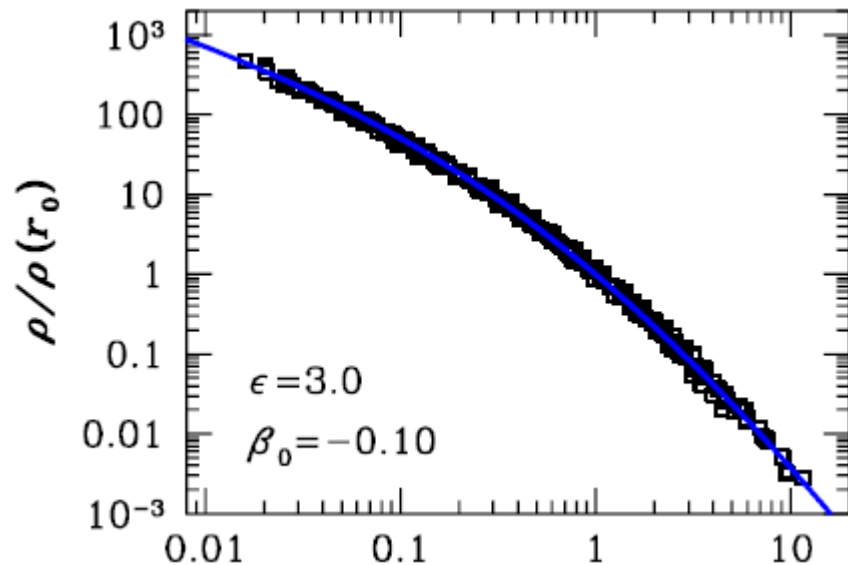
Hansen & Moore 2006

Dark matter halo models

Dehnen & MacLaughlin (2005): Jeans models for dark halos based on these assumptions

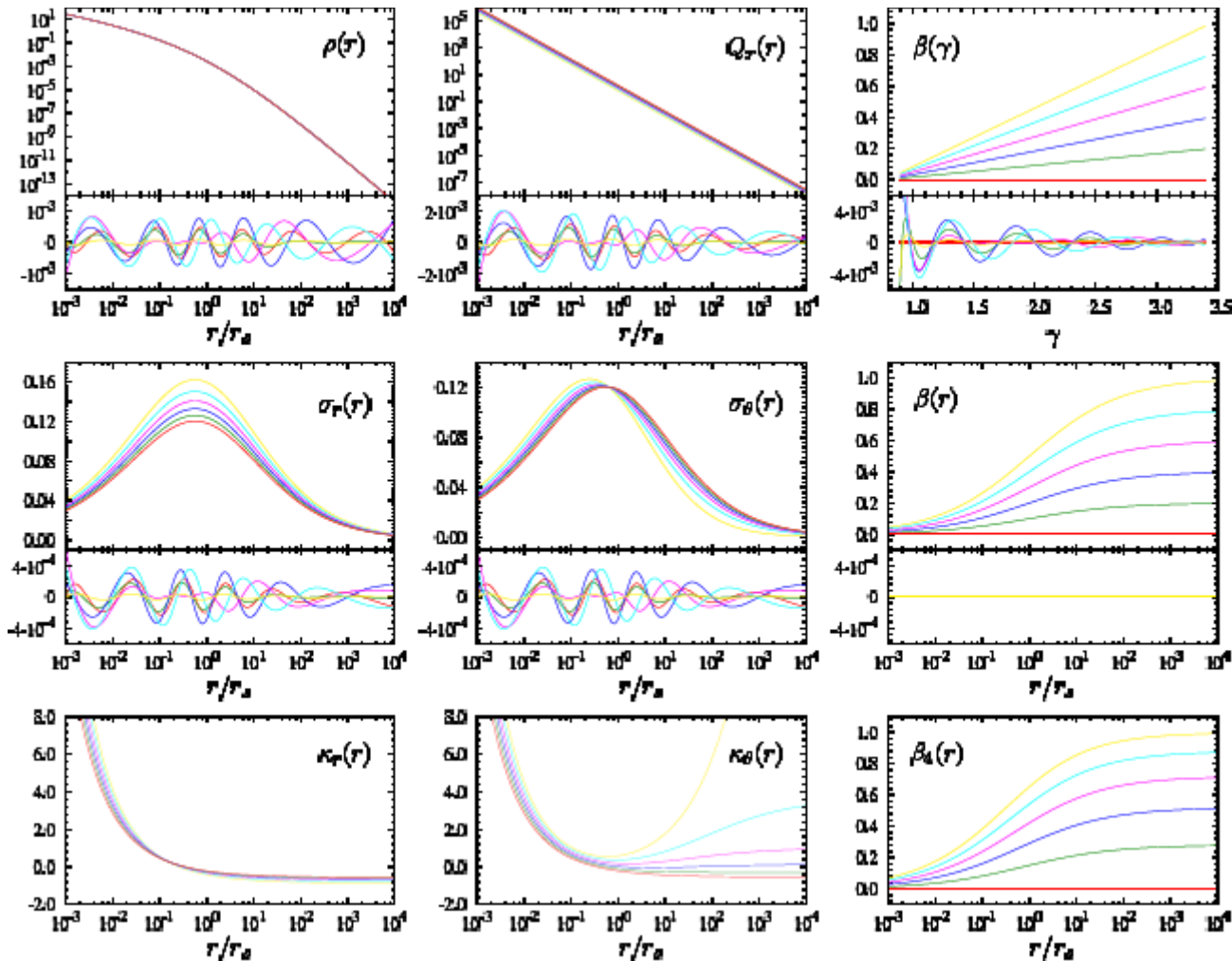
- potential, density and dispersion are analytical
- anisotropy profile as in our components
- excellent fit to a set of simulated halos

$$\gamma(r) = \frac{\gamma_0 + \gamma_\infty x^\eta}{1 + x^\eta} \quad \frac{\rho}{\sigma_r^\epsilon}(r) = \frac{\rho}{\sigma_r^\epsilon}(r_s) \left(\frac{r}{r_s} \right)^{-\alpha_{\text{crit}}} \quad \beta(r) = \frac{\beta_0 + \beta_\infty x^\eta}{1 + x^\eta}$$



Dark matter halo models

- dynamical models with $N=10$ components
- models with a range of outer anisotropies from 0 to 1
- quality of the fits is excellent

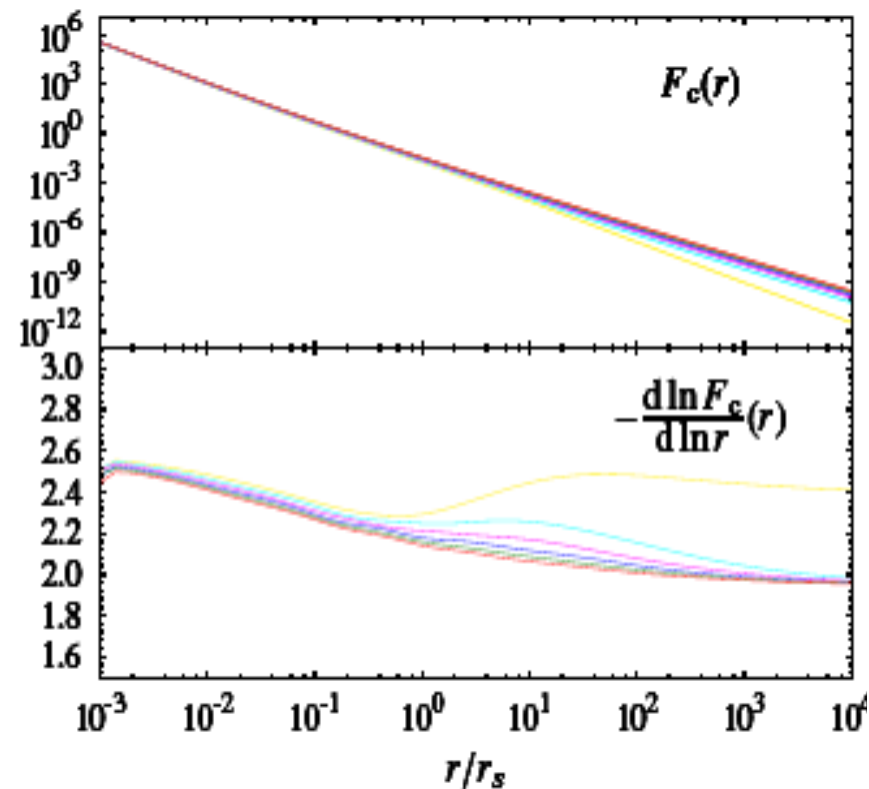
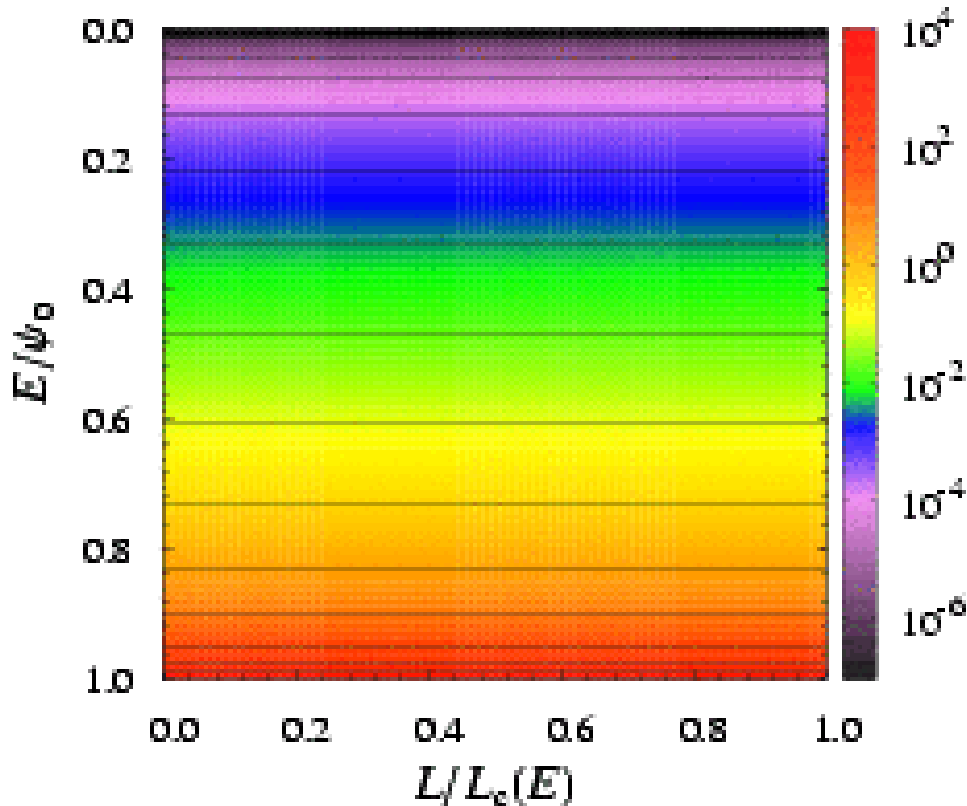


Van Hese, Baes
& Dejonghe
2008, ApJ

Distribution function

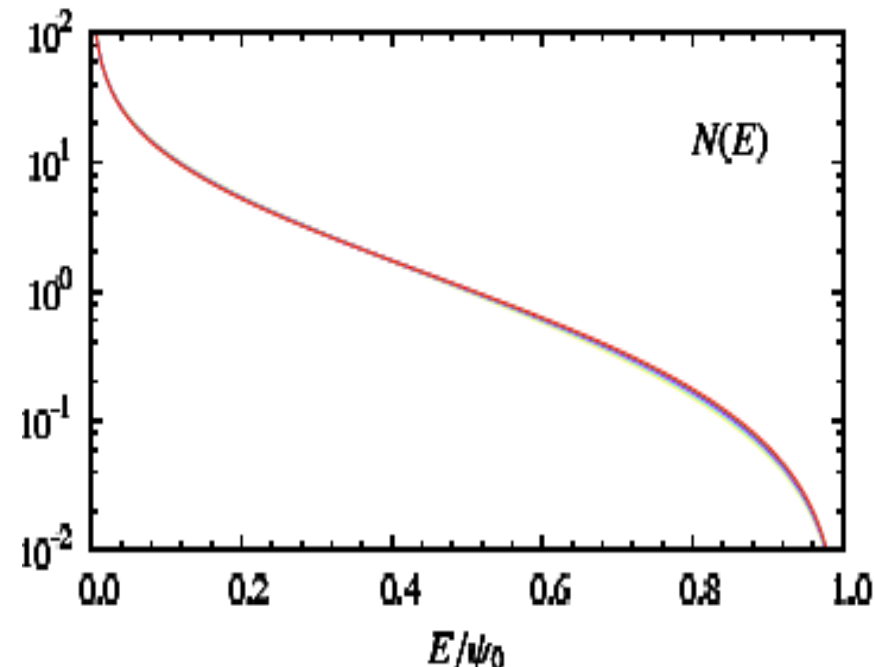
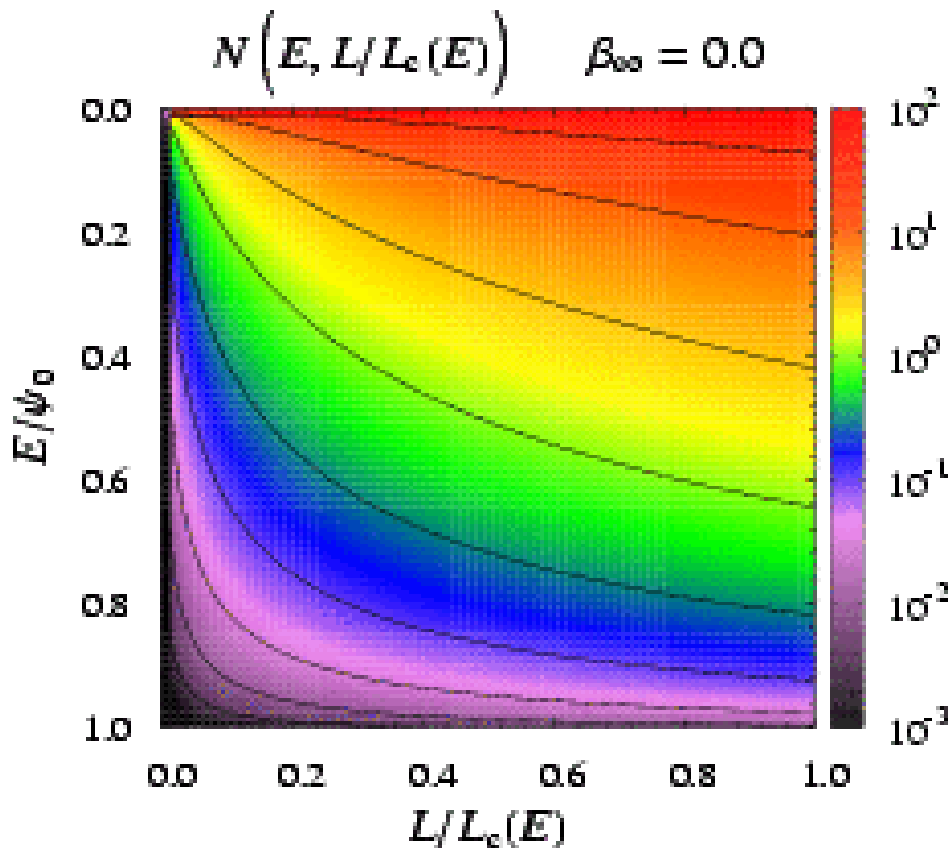
- DF positive for all anisotropies
- smoothly varying for β_{∞} ranging from 0 to 1
- remarkable: DF for circular orbits is power-law

$F(E, L)$ $\beta_{\infty} = 0.0$



Differential distribution

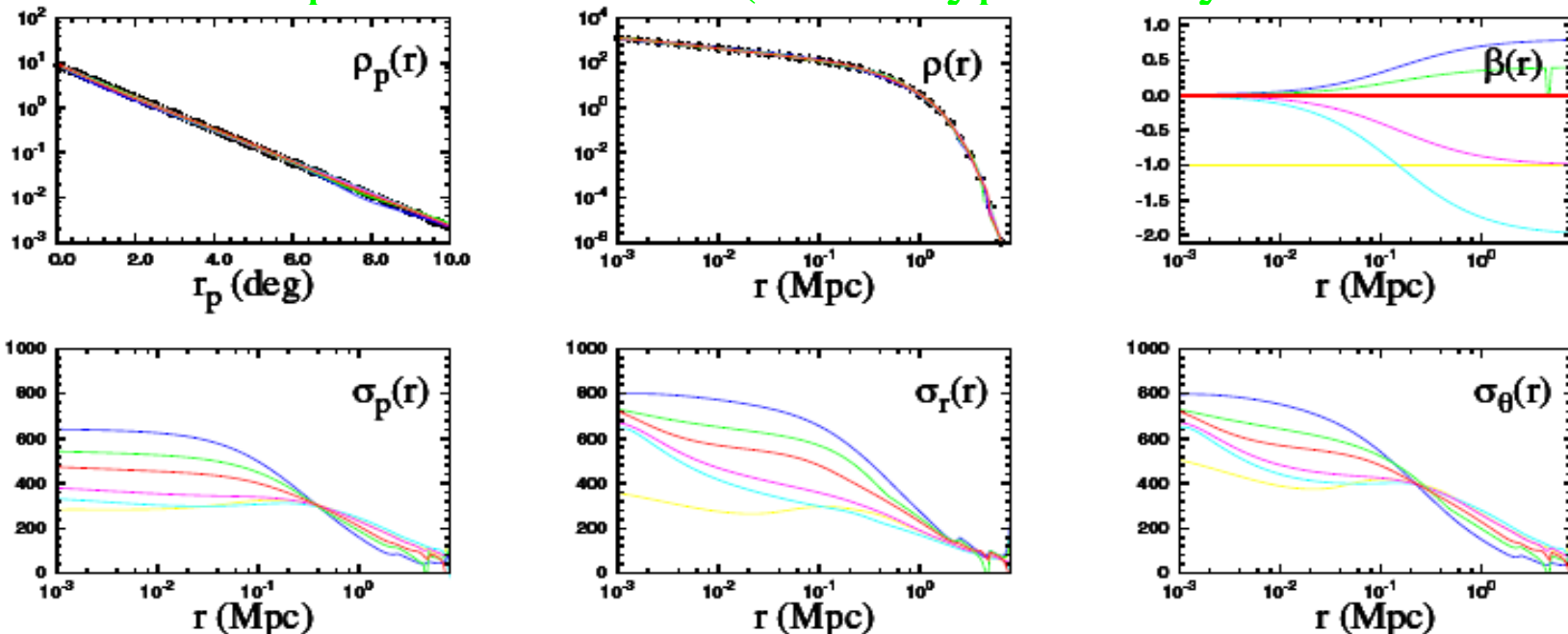
- $N(E, L)$ gives the number of orbits with integrals E and L
- $N(E)$ nearly independent of the anisotropy



Future work

- Modeling different potentials and densities
- Gain more insight into the dynamical structure
- Check other “universal” properties (e.g. velocity distribution)
- Fitting directly to N -body data
- Modeling dark matter and galaxy distribution in clusters
- Extension to non-spherical halos

Dwarf Ellipticals in Fornax (Sersic-type density in NFW-halo)



Conclusions

- new technique to construct analytical dynamical models with arbitrary density and realistic anisotropy profile
- useful in other dynamical studies, with other data moments than the density
- our models can generate initial conditions for N -body simulations
- phase-space investigation of dark matter halos with universal properties
 - very accurate fits
 - Jeans models of DM05 have positive DF
 - $N(E)$ nearly independent of anisotropy

valuable complementary approach to N -body studies