

Dark Matter Annihilation in the Late Universe

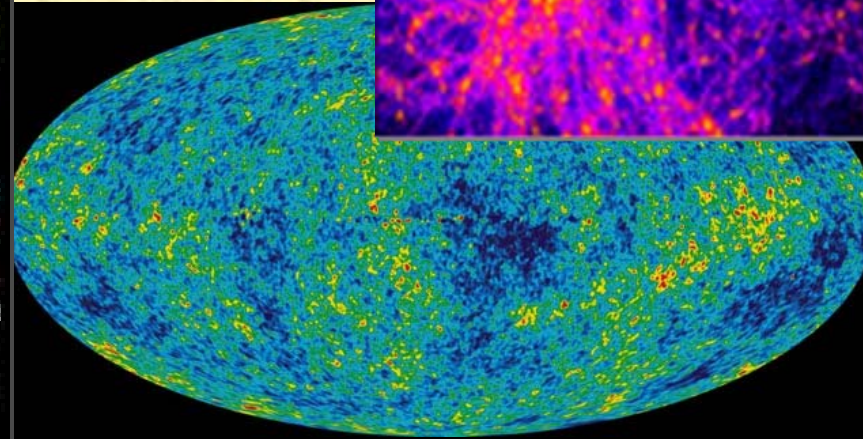
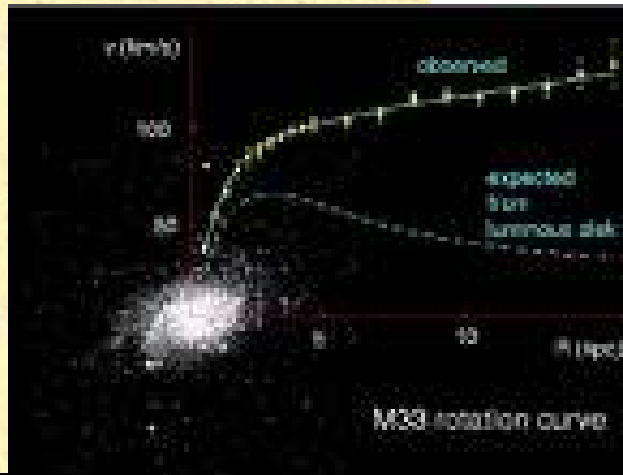
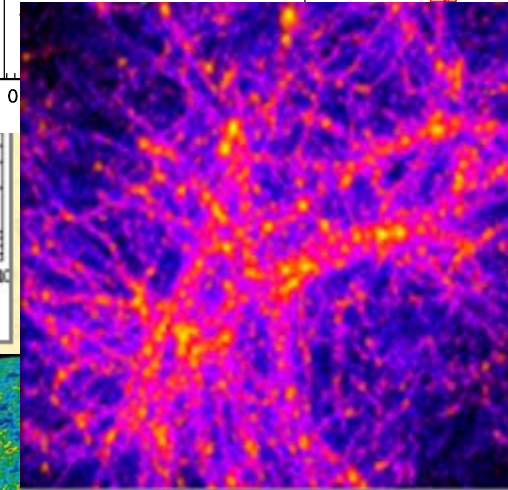
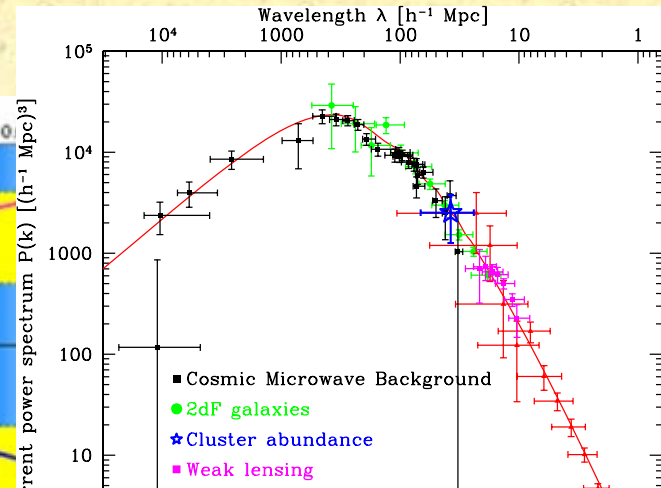
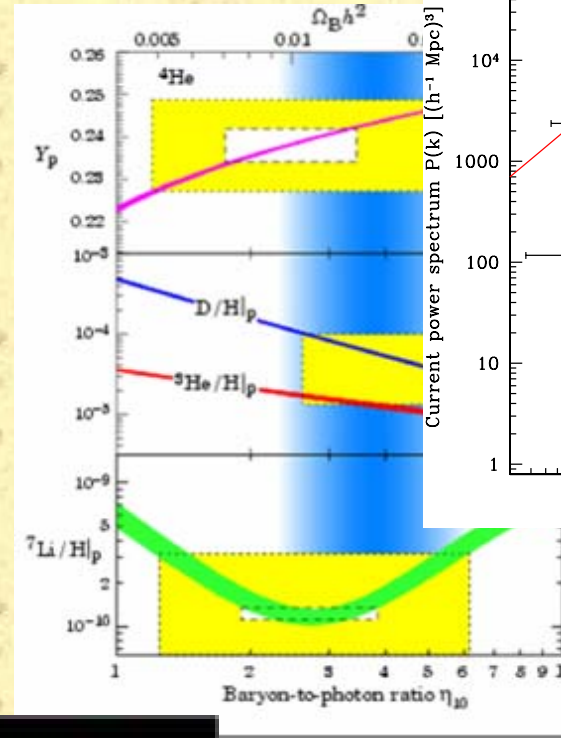
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John Beacom (Ohio State), Gregory Mack (Ohio State), Hasan Yüksel (Ohio State), Thomas Jacques (Melbourne), Gianfranco Bertone (Paris), James Dent (Vanderbilt) and Tom Weiler (Vanderbilt).

Dark Matter

- Evidence for dark matter arises from a wide range of astrophysical observations which are each sensitive to dark matter's *gravitational influence*.
- As yet, very little information on the *particle properties* of dark matter.



Dark Matter Candidate Zoo

Axions, Neutralinos, Gravitinos, Axinos, Sneutrinos, Kaluza-Klein particles, Heavy Fourth Generation Neutrinos, Mirror particles, superWIMPs, WIMPzillas, Sterile Neutrinos, Light Scalars, Q-Balls, Brane World Dark Matter, Primordial Black Holes, things we haven't thought of

It's clear that any model-dependent approach to searching for dark matter has shortcomings.

Dark Matter Annihilation

- ❖ Self-annihilation cross-section is a fundamental property of dark matter
- ❖ For thermal relics it is specified by the DM relic density:
e.g. $\Omega_{DM} \sim 0.3$ implies $\langle \sigma_A v \rangle \sim 3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$
- ❖ More generally, it controls the DM annihilation rate in galaxies today \rightarrow can affect galaxy halo density profiles

There exist two general constraints on the rate of dark matter annihilation/*disappearance*:

Constraint 1. - Unitarity Bound

❖ Unitarity sets a general upper bound on the cross-section:

$$\langle \sigma_A v \rangle \leq 4\pi / m_\chi^2 v$$

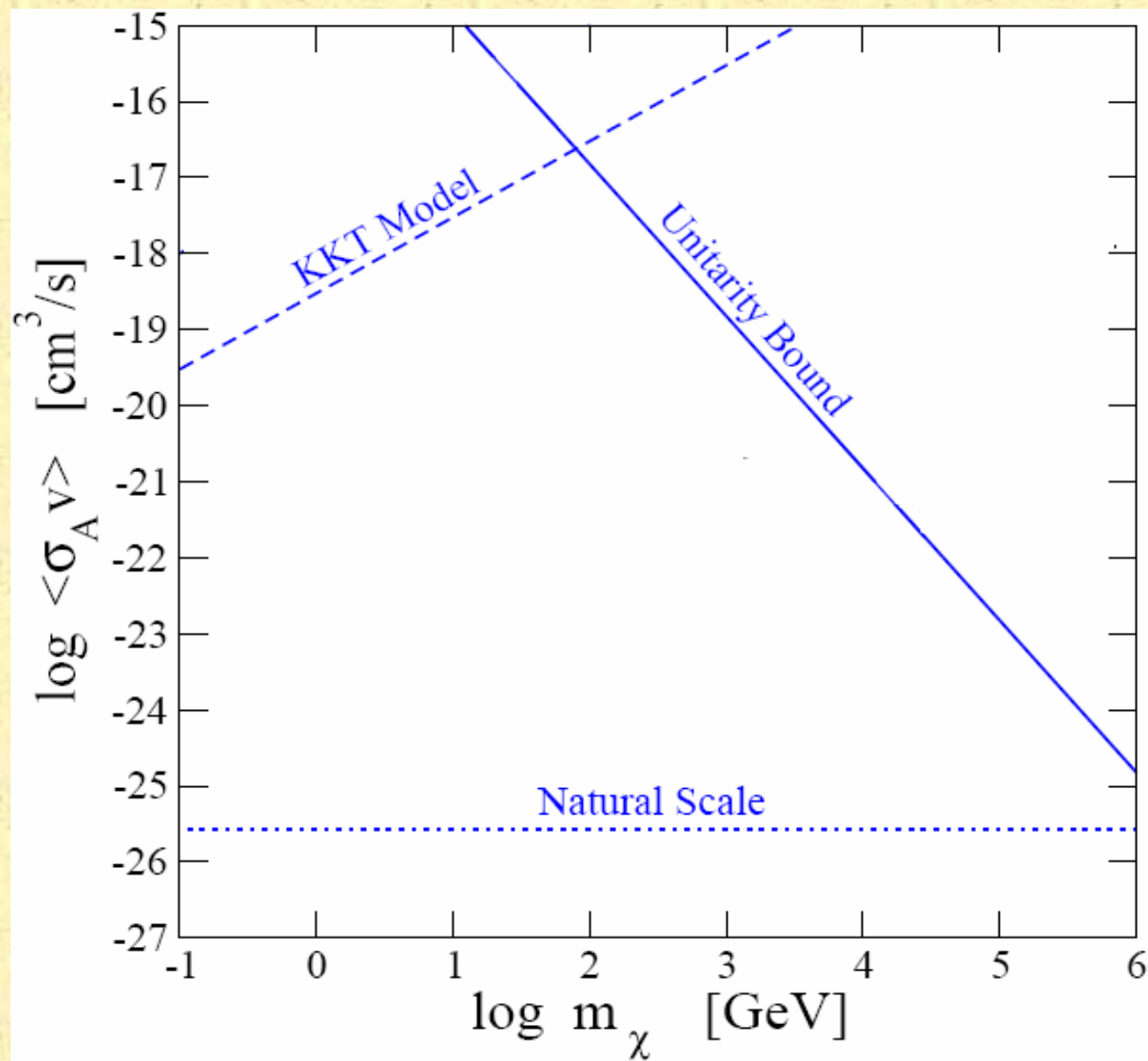
(in the low-velocity limit, where the cross-section is s-wave dominated)

→ In galaxies today:

$$\langle \sigma_A v \rangle \leq 1.5 \times 10^{-13} \frac{\text{cm}^3}{\text{s}} \left[\frac{\text{GeV}}{m_\chi} \right]^2 \left[\frac{300 \text{ km/s}}{v_{rms}} \right].$$

L. Hui

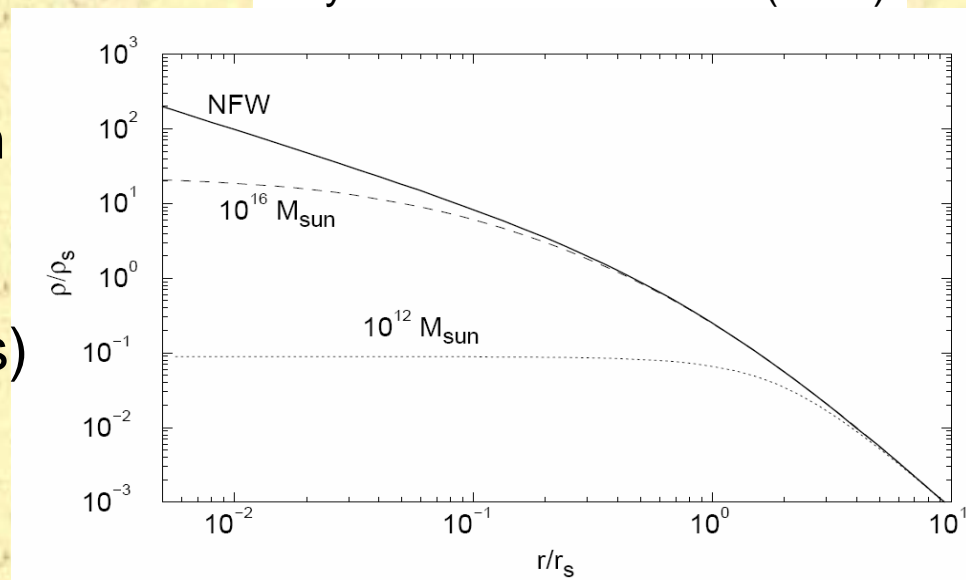
❖ Most restrictive for high masses.



Constraint 2. - Kaplinghat-Knox-Turner Model

Phys. Rev. Lett. 85. 3335 (2000)

❖ Large dark matter annihilation rate flattens galaxy cores
 → invoked to resolve conflict between predicted (sharp cusps) and observed (flat cores) halo density profiles.



❖ KKT model requires cross-sections $\sim 10^7$ times larger than the natural scale for a thermal relic:

$$\langle \sigma_A v \rangle_{\text{KKT}} \simeq 3 \times 10^{-19} \frac{\text{cm}^3}{\text{s}} \left[\frac{m_\chi}{\text{GeV}} \right]$$

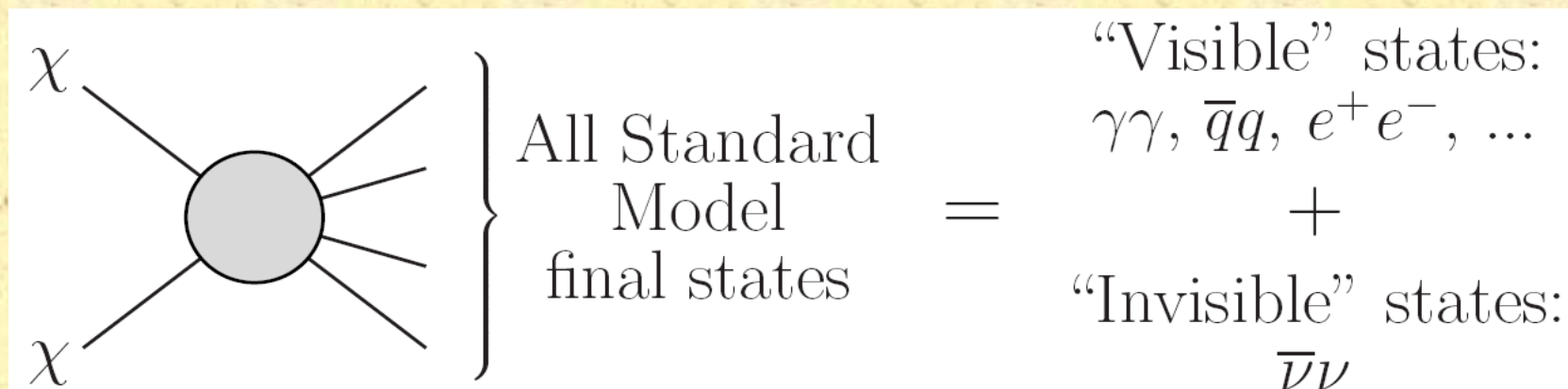
❖ Reinterpret this type of model as upper bound on $\langle \sigma_A v \rangle$

Dark Matter Annihilation & Neutrinos → Indirect Detection

Search for fluxes of DM annihilation or decay products:

- ❖ Annihilation in our Galaxy
→ look for flux coming from Galactic center
- ❖ Annihilations in galaxies throughout the universe
→ cosmic diffuse flux
- ❖ Despite being harder to detect than gamma rays, neutrinos provide important information and ***strong, model independent bounds.***

If the dark matter is the lightest new particle:



❖ All final states except neutrinos produce gamma rays,

→ Bound the *total* cross-section with the neutrino signal limit
 i.e. Assume $\text{Br}(\text{“invisible”}) = 100\%$

Diffuse Neutrino Signal

❖ Annihilation to neutrinos:

$$\chi\chi \rightarrow \bar{\nu}\nu$$

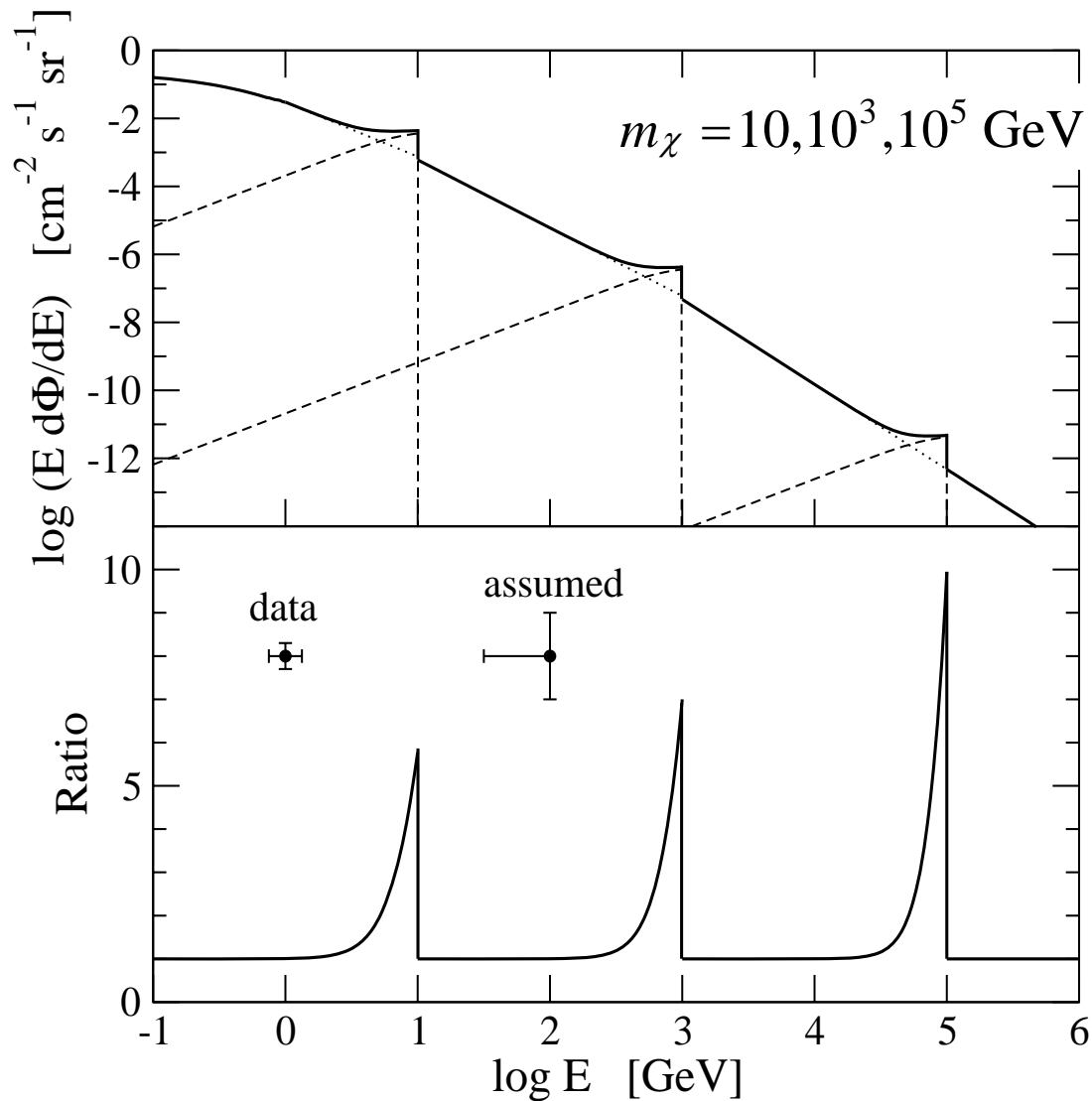
❖ Diffuse flux (Ullio, et al.):

$$\frac{d\Phi_\nu}{dE} = \frac{\langle\sigma_A v\rangle}{2} \frac{c}{4\pi H_0} \frac{\Omega_{DM}^2 \rho_{\text{crit}}^2}{m_\chi^2} \int_0^{z_{\text{up}}} dz \frac{\Delta^2(z)}{h(z)} \frac{dN_\nu(E')}{dE'}$$

The factor Δ accounts for the increase in density due to the clustering of dark matter in halos. ($\Delta=1$ corresponds to all matter being at the average density in the Universe today)

❖ Atmospheric neutrinos are the background

❖ Conservative detection criteria: Signal 100% as large as angle averaged atmospheric neutrino background.



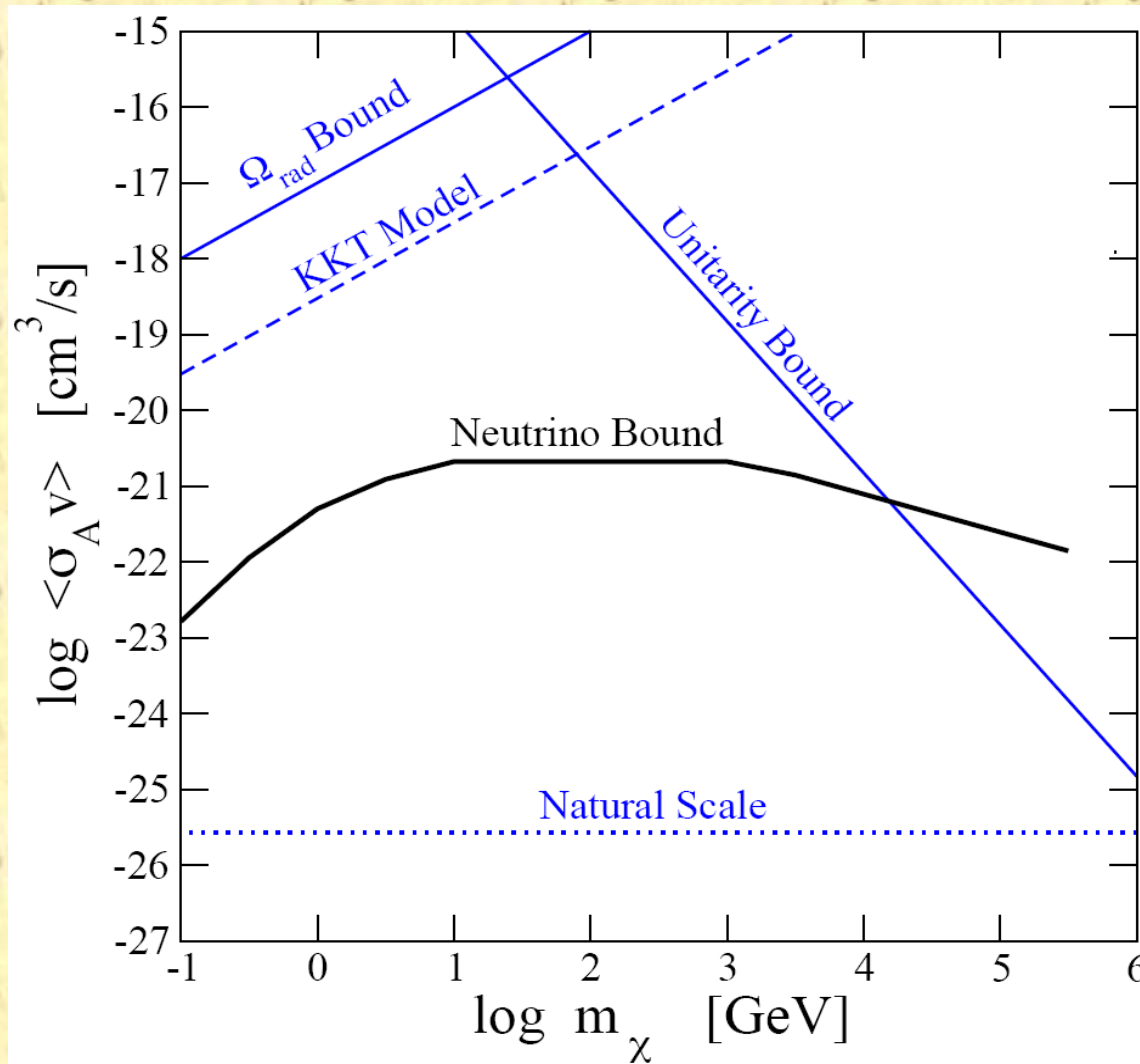
Upper panel: –

Annihilation flux
superimposed on
atmospheric neutrino
background

Lower panel:

$(\text{Signal} + \text{Background}) /$
 Background

Upper bounds on the dark matter total annihilation cross-section



Beacom, Bell, & Mack,
PRL99, 231301, 2007.

Also:

Yuksel, Horiuchi,
Beacom, & Ando,
PRD 76, 123506, 2007.

Kachelriess and Serpico,
PRD 76, 063516 (2007).

Bell, Dent, Jacques, &
Weiler, arXiv:0805.3423

Annihilation effect in Galaxy Halos?

- ❖ Annihilation flattens halo cusps to a core density of:

$$\rho_A \sim m_\chi / (\langle \sigma_A v \rangle H_0^{-1})$$

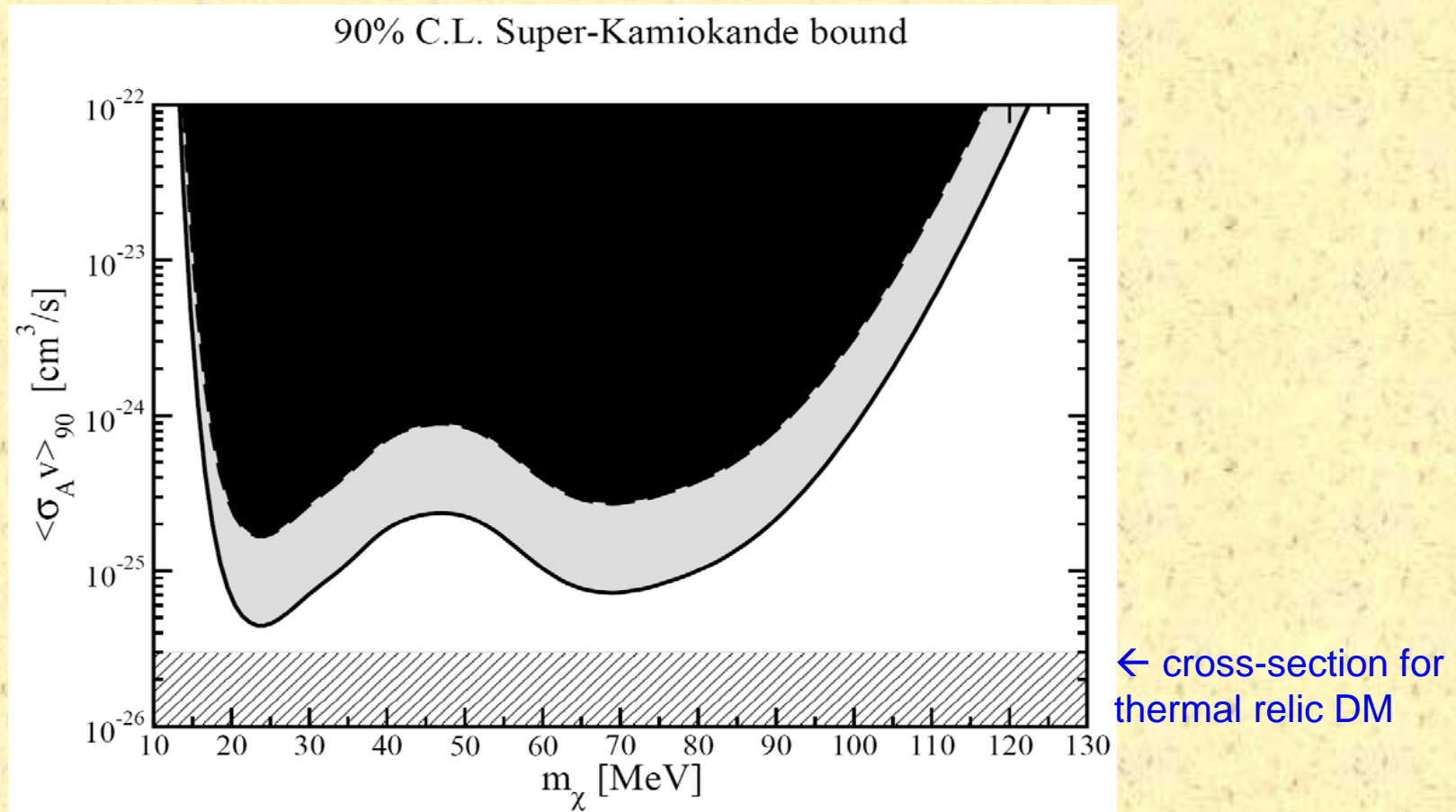
- ❖ Our bound implies that for all $m > 0.1$ GeV:

$$\rho_A \gtrsim 5 \times 10^3 \text{ GeV/cm}^3$$

→ Only affects the very inner region of typical galaxies.
e.g. In the Milky Way, this density occurs only at radii < 1 pc for an NFW profile (and maybe not at all for less steep profiles).

Dark matter annihilation cannot have a macroscopic effect on galactic halos.

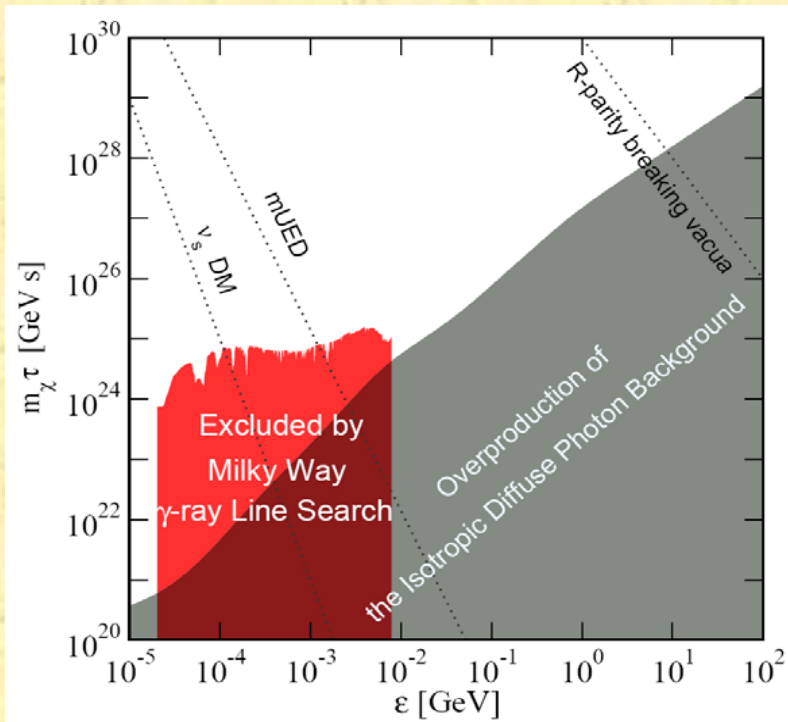
Dark matter annihilation – MeV mass



Palomares-Ruiz & Pascoli, PRD 77, 025025 (2008)

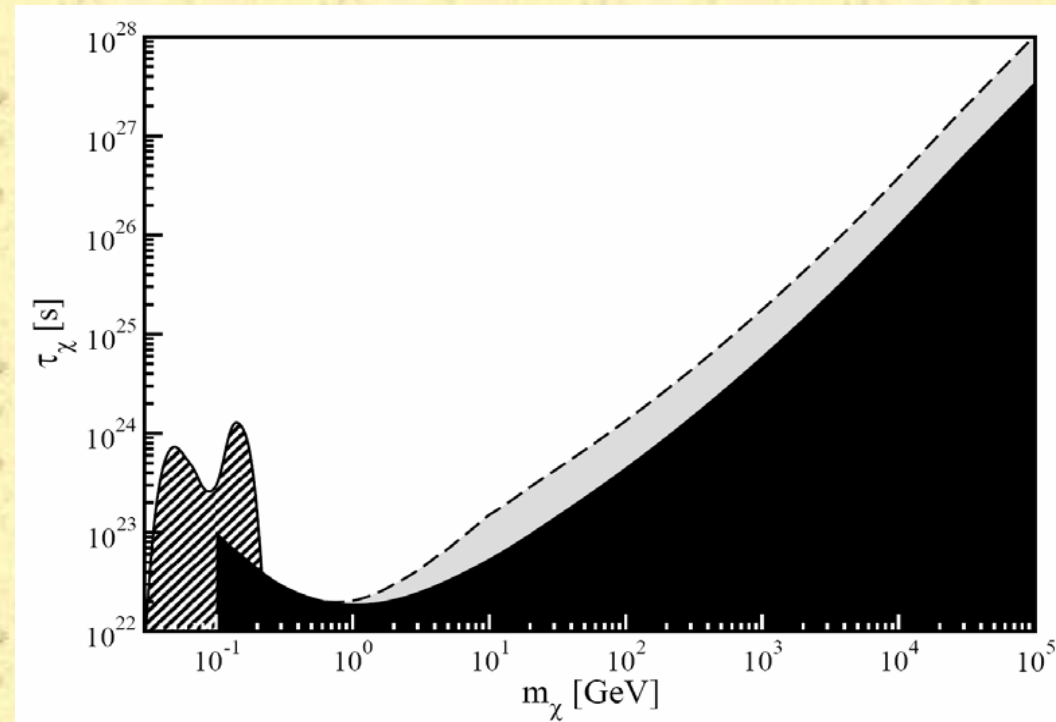
Dark Matter Lifetime

Photons $X \rightarrow X' \gamma$
... strong limit



Yuksel & Kistler

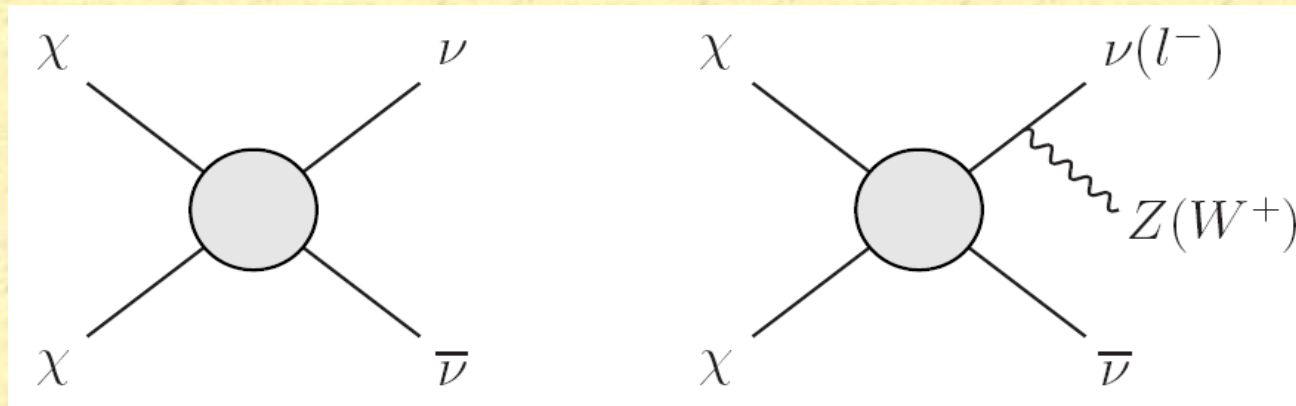
Neutrinos $X \rightarrow \nu \nu$
... model-independent limit



Palomares-Ruiz

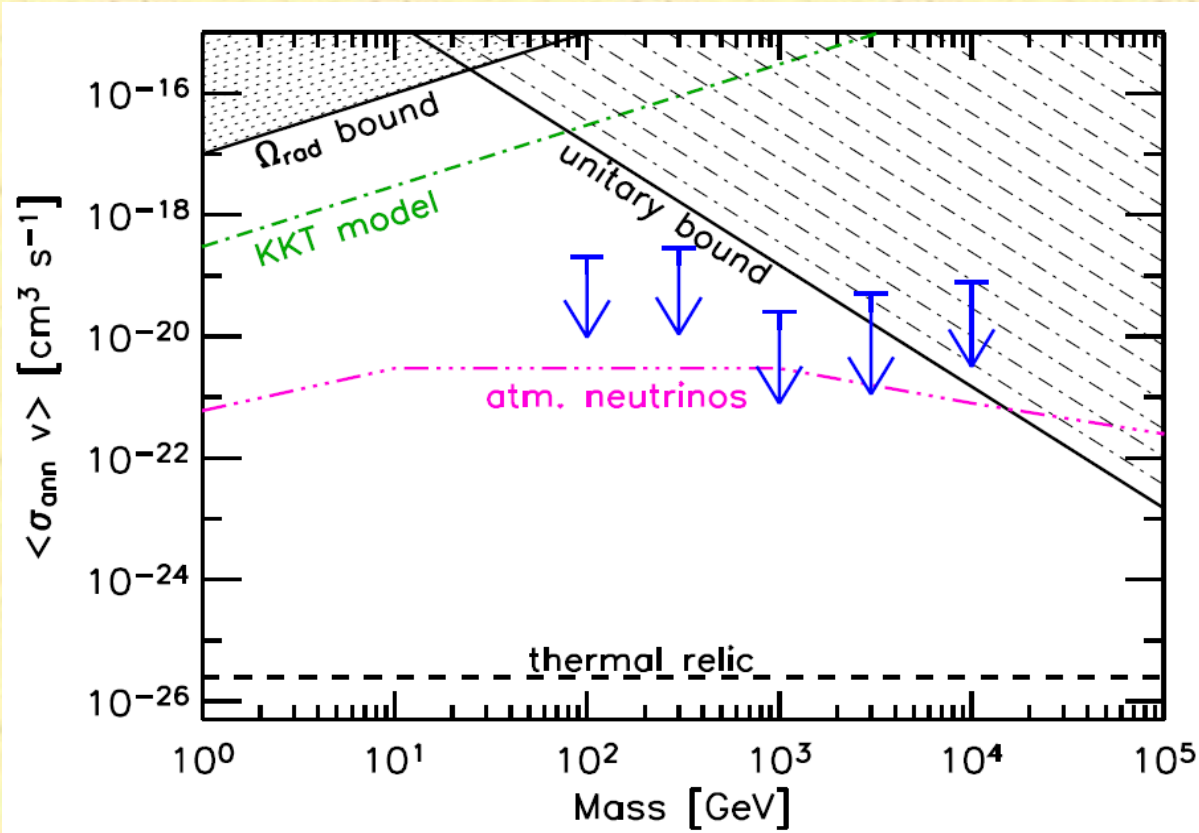
Radiative corrections \rightarrow photons

$\chi\chi \rightarrow \nu\bar{\nu}$ is accompanied by: $\chi\chi \rightarrow \nu\bar{\nu}Z \rightarrow$ gamma rays



- ❖ Z decays to hadrons $\sim 70\%$ of the time.
- ❖ On average, 20 neutral pions per Z decay, which in turn decay as $\pi \rightarrow \gamma\gamma$
- ❖ For $E \gg M_Z$: ratio of the cross-sections $\sim \alpha$

Radiative corrections to neutrino processes give photons

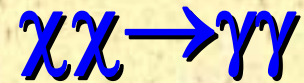


But direct
neutrino
limits are
comparable/
stronger!

Kachelriess and Serpico,
PRD 76, 063516 (2007).

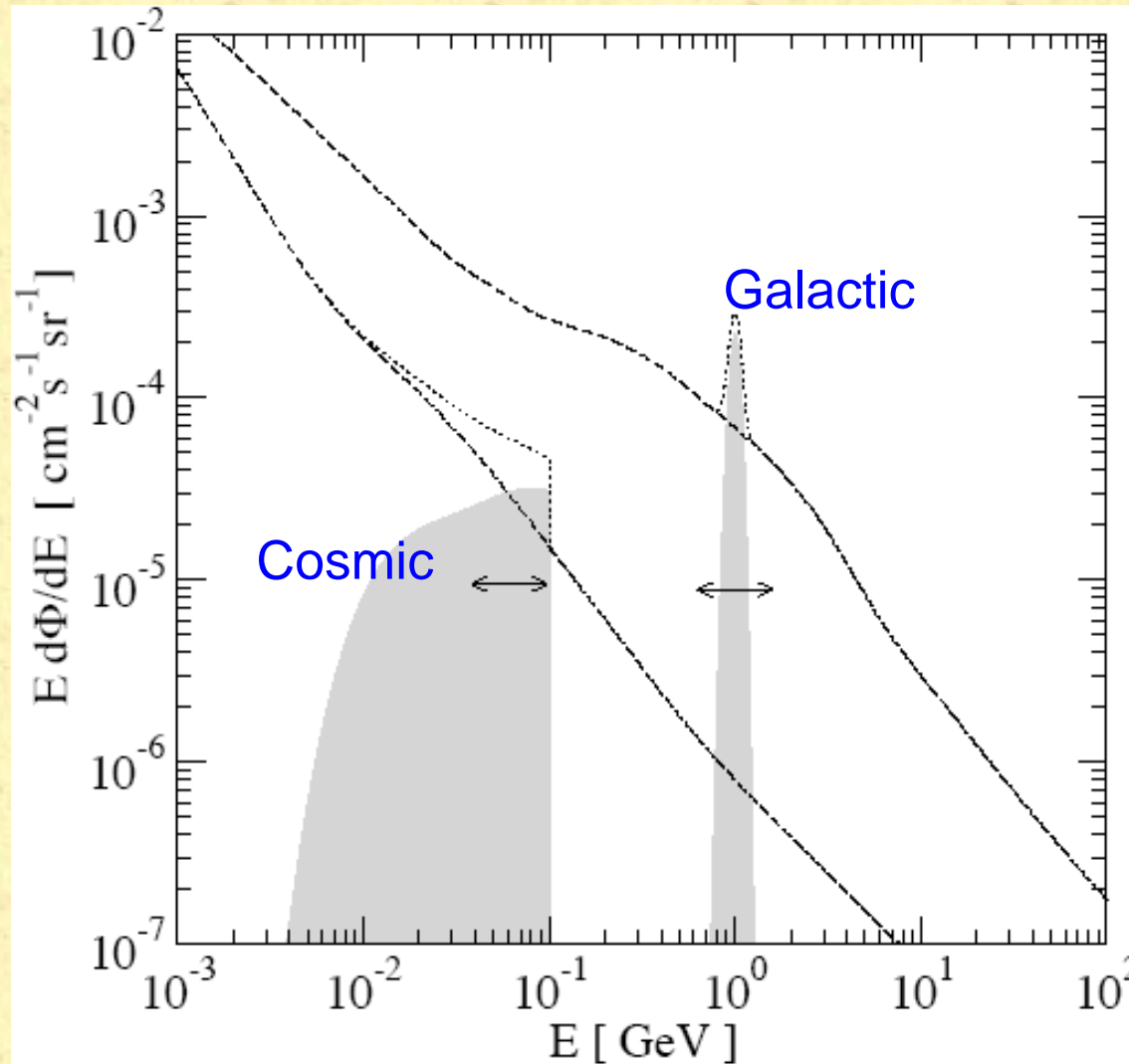
Bell, Dent, Jacques, &
Weiler, arXiv:0805.3423

Direct Annihilation to gamma rays



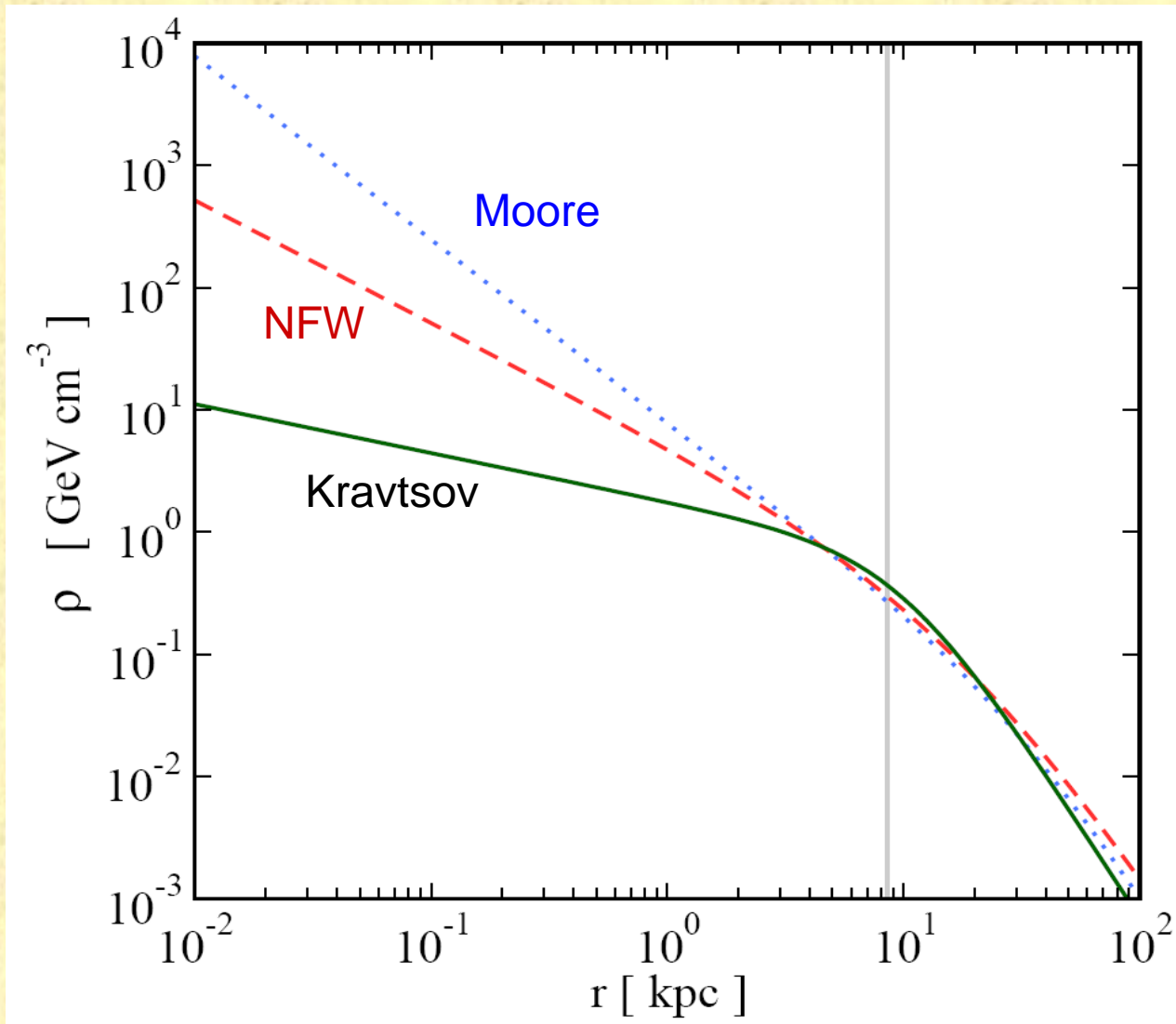
- ❖ The process $\chi\chi \rightarrow \gamma\gamma$ is present in most dark matter models, though with a model-dependent rate. (Usually occurs at loop order.)
- ❖ It's detection would be a “smoking gun” signature of dark matter.
- ❖ Monochromatic line at $E_\gamma = M_\chi$ would measure DM mass.

Annihilation signals & γ -ray background



Mack, Jacques,
Beacom, Bell,
Yuksel
arXiv:0803.0157

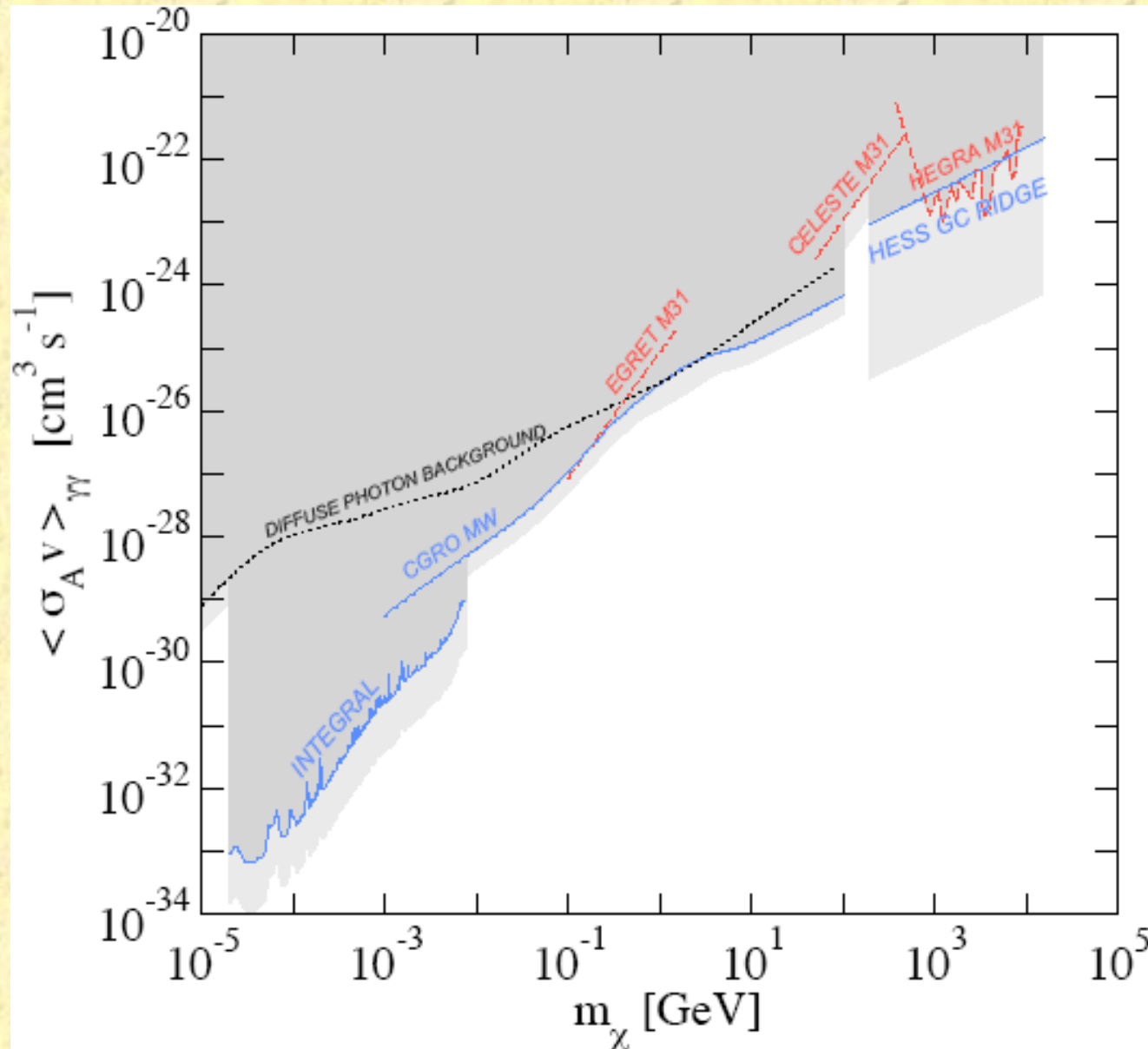
DM halo density vs radius



Yuksel, Horiuchi,
Beacom, & Ando,
PRD 76, 123506,
2007.

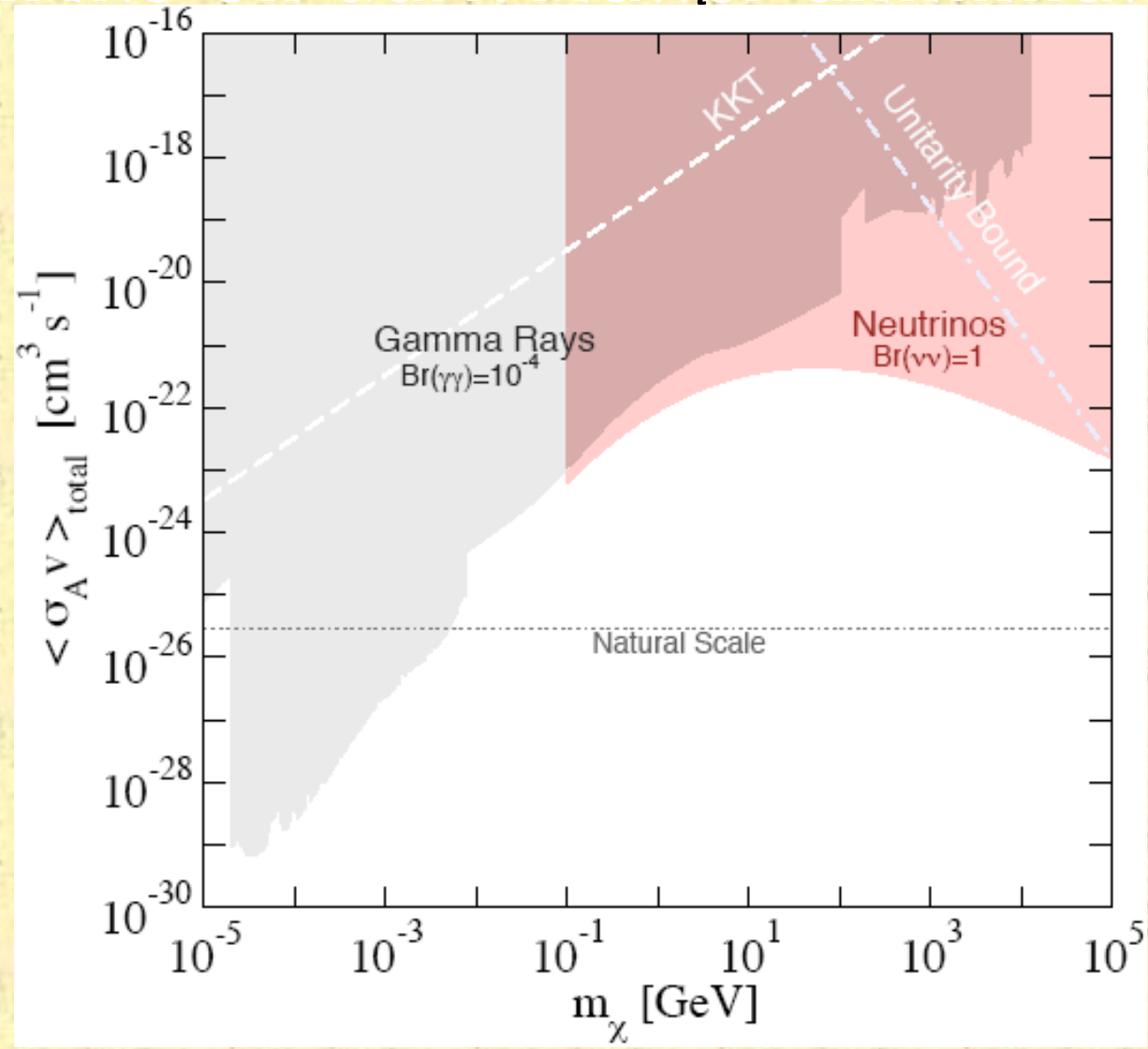
Cross-section limits for $\chi\chi \rightarrow \gamma\gamma$

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Mack, Jacques,
Beacom, Bell,
Yuksel
arXiv:0803.0157

Comparison of neutrino and photon limits on dark matter annihilation



Mack, Jacques,
Beacom, Bell,
Yuksel
arXiv:0803.0157

Summary – DM Annihilation

- ❖ Dark matter total annihilation cross-section can be bounded using the least detectable annihilation products (i.e. neutrino appearance rate.)
- ❖ Neutrino bound much stronger than Unitarity for $m < 10$ TeV.
- ❖ For large mass, neutrinos bounds are comparable to those for $\gamma\gamma$
- ❖ Dark Matter halos cannot be significantly modified by annihilation.